

Strong decays of excited baryons in Large N_c QCDJ.L. Goity^{1,2} and N.N. Scoccola^{3,4,5}¹ Physics Dept., Hampton University, Hampton, VA 23668, USA.² TJNAF, Newport News, VA 23606, USA.³ CONICET, Rivadavia 1917, 1033 Buenos Aires, Argentina.⁴ Lab. TANDAR, CNEA, Av.Libertador 8250, 1429 Buenos Aires, Argentina.⁵ Universidad Favaloro, Solís 453, 1078 Buenos Aires, Argentina.*Received 31 October 2005*

Abstract. We present the analysis of the strong decays widths of excited baryons in the framework of the $1/N_c$ expansion of QCD. These studies are performed up to order $1/N_c$ and include both positive and negative parity excited baryons.

Keywords: Excited baryons, strong decays, large N_c QCD

PACS: 14.20.Gk, 14.20.Jn, 12.39.Jh, 12.40.Yx

We consider here some application of the $1/N_c$ expansion of QCD[1,2] to the study of the strong decays of excited baryons belonging to the **70**-plet of negative parity and the **56**-plet of positive parity. For simplicity we consider here only the non-strange baryons of those multiplets. The ℓ_P partial wave decay width into a ground state (GS) baryon and a pseudo-scalar meson with isospin I_P is given by

$$\Gamma^{[\ell_P, I_P]} = \frac{k_P}{8\pi^2} \frac{M_{B^*}}{M_B} \frac{|B(\ell_P, I_P, S, I, J^*, I^*, S^*)|^2}{\sqrt{(2J^* + 1)(2I^* + 1)}}, \quad (1)$$

where $B(\ell_P, I_P, S, I, J^*, I^*, S^*)$ are the reduced matrix elements of the baryonic operator. Such operator admits an expansion in $1/N_c$ that has the general form:

$$B_{[\mu, \alpha]}^{[\ell_P, I_P]} = \left(\frac{k_P}{\Lambda}\right)^{\ell_P} \sum_q C_q^{[\ell_P, I_P]}(k_P) \sum_m \langle \ell, m; j, j_z | \ell_P, \mu \rangle \xi_m^\ell \left(\mathcal{G}_{[j_z, \alpha]}^{[j, I_P]} \right)_q, \quad (2)$$

The factor $\left(\frac{k_P}{\Lambda}\right)^{\ell_P}$ is included to take into account the chief meson momentum dependence of the partial wave. For definiteness, the scale Λ is taken to be equal to 200 MeV. The operator ξ_m^ℓ drives the transition from the $(2\ell + 1)$ -plet to the singlet $O(3)$ state, and the spin-flavor operators $\left(\mathcal{G}_{[j_z, \alpha]}^{[j, I_P]} \right)_q$ give transitions within the

spin-flavor representations in which the excited and GS baryons reside. The label j denotes the spin of the spin-flavor operator while its isospin obviously coincides with the isospin of the emitted meson. The dynamics of the decays is encoded in the effective dimensionless coefficients $C_q^{[\ell_P, I_P]}(k_P)$. The relevant list of independent operators \mathcal{G} for the decays we are interested as well as the corresponding fits and calculated widths in can be found in Refs.[3,4]. From those results we observe that, in general, the LO operators already provide a reasonable description of the decay widths. Exception to this are the D-wave decays of the $(1^-, \mathbf{70})$ -plet where the NLO corrections play some role in improving the results. Another interesting observation that points to the consistency of the framework based on the approximate $O(3) \times SU(2N_f)$ symmetry is that the predicted suppression of the η channels in the decays positive parity baryons is clearly displayed by the observed decays. This represents a strong experimental confirmation that those excited baryons do belong primarily into a symmetric representation of $SU(2N_f)$. In the case of the negative parity baryons the η channel is not suppressed and is indeed very important. This implies that these states belong primarily into the mixed-symmetric representation of $SU(2N_f)$. Thus, η channels serve as a selector of the spin-flavor structure of excited baryons.

To conclude, the $1/N_c$ expansion provides a systematic approach to the properties of the excited baryons. In the case of the negative parity $(1^-, \mathbf{70})$ -plet and the positive parity $(2^+, \mathbf{56})$ baryons it successfully describes the existing data and, to the order considered, also makes numerous testable predictions.

Acknowledgments

This work was partially supported by the National Science Foundation (USA) through grant # PHY-9733343 (JLG) and by the CONICET and ANPCYT (Argentina) through grants PIP 02368 and PICT 03-08580 (NNS).

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