

# Experimental study of Bernoulli's equation with losses

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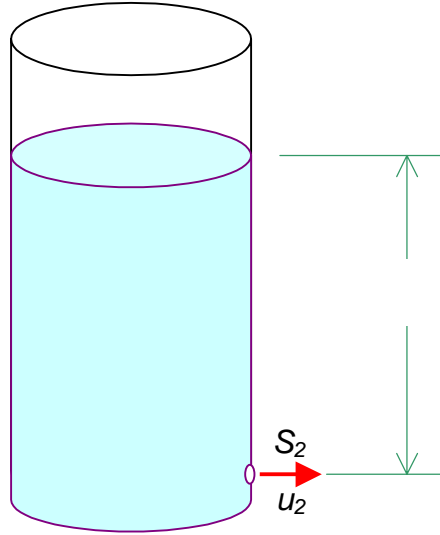
The purpose of this work is to present a simple and inexpensive experiment to study the drainage of a cylindrical vessel. The experimental set up consists of a transparent cylinder and a web cam or a digital camera connected to a computer. The model proposed to explain the results makes use of Bernoulli's equation for real flows that includes energy losses. The experiment allows us to test thoroughly the implication of the model and to extract the relevant parameter associated with the energy losses. The experimental results can be nicely explained by the model, which is a generalization of Torricelli's expression, and clearly illustrate the utility of the extended Bernoulli equation for real fluids.

## I. – Introduction

The Bernoulli equation including losses is an useful and important expression of wide application in many branches of science and engineering.<sup>1,2,3</sup> General forms of Bernoulli's equation, also known as the *extended Bernoulli's equation*, valid for viscous fluids, have been discussed previously in this journal<sup>4</sup>. Nonetheless, there are relatively few experiments,<sup>5,6</sup> accessible to beginner and intermediate students that illustrate its use and applications. In this work we present a conceptually simple and inexpensive experiment to study the drainage of a cylindrical vessel that, we believe, clearly illustrates an application of Bernoulli's equation for real liquids. It is essentially a recreation of Torricelli's experiment, profiting from the advantages of new technologies. The experimental set up consists essentially of a transparent cylinder and a web cam connected to a computer. In the first part of this study we present a model based on the Bernoulli equation for real flows. Then we present the basic characteristic of the experiment and the experimental results. The experiment allows us to thoroughly test the implications of the model and to extract the relevant parameter associated with the energy losses. Within this context, we find this experiment pedagogically useful to introduce the Bernoulli equation for real flows to beginner students. The experiment involves concepts that are relatively simple to discuss theoretically, the physics is easy to visualize and it is quite straightforward to quantitatively test the implication of the model.

## II. Theoretical considerations

For Newtonian fluids the shear stress is proportional to the velocity gradient, therefore the velocity of the fluid on the surface of a solid must be zero, otherwise the velocity gradient and the shear stress will be infinite. Only for an ideal fluid, that is a fluid with zero viscosity,  $\eta=0$ , is it possible to have finite velocity on the surface of a solid. Therefore, when a real fluid flows through the interior of a tube or between two surfaces, there are two effects that are a consequence of the viscosity. The velocity profile presents a maximum at the center of the tube. The other important consequence of the viscosity is the nonconservation of mechanical energy in the system or the presence of energy losses.



**Figure 1:** Schematic diagram of the experimental setup

In its simplest form, the Bernoulli equation is basically a statement of the conservation of mechanical energy per unit of volume along a stream line.<sup>1,2</sup> In the presence of viscosity, Bernoulli's equation becomes an expression of the energy balance, often expressed in terms of energy per unit of volume or pressure (or energy per unit of weight or head) between two points in the flow of fluid. The Bernoulli equation for a steady flow of real fluids in a pipe can be written in the following form:<sup>1,2,3,4</sup>

$$\frac{1}{2} \alpha_1 u_1^2 + \frac{P_1}{\rho_1} + g z_1 = \frac{1}{2} \alpha_2 u_2^2 + \frac{P_2}{\rho_2} + g z_2 + \Delta w_{loss} . \quad (1)$$

Indexes 1 and 2 refer to two different points in the flow, 1 being upstream of 2,  $g$  represents the local acceleration of gravity,  $P$  the pressure,  $z$  the vertical height of the point, and  $u$  the average velocity of the flow along the tube. If the fluid is incompressible, then  $\rho_1 = \rho_2$ . The average velocity is defined in terms of the flux  $Q$  as:

$$Q = \iint_S \vec{v} \cdot d\vec{S} = u \cdot S . \quad (2)$$

Where  $S$  is the area of the normal cross section of the flow (or pipe) and  $v$  is the local velocity. Expression (2) can be regarded as the definition of the average velocity  $u$ . The *kinetic energy coefficients*  $\alpha_i$  in Eq. (1) represent the ratio between actual kinetic energy that flows through a

normal cross section in the flow and the kinetic energy of the same flux, but with uniform velocity profile equal to  $u$ , more specifically  $\alpha$  is defined as<sup>1,3</sup>

$$\alpha = \frac{\iint_S v^3 \cdot dS}{u^3 \cdot S}. \quad (3)$$

Only for a uniform profile  $\alpha$  is equal to 1, for any non uniform profile  $\alpha > 1$ . For a parabolic profile of velocity (laminar flow)  $\alpha = 2$ . The term  $\Delta w_{loss}$  in expression (1) represents the head loss that is proportional to the mechanical energy loss between the points 1 and 2. For a great variety of actual situations,  $\Delta w_{loss}$  can be expressed as the sum of two terms, one dependent of the square of the average velocity and one independent of the velocity. The term proportional to  $u^2$  is often referred to as minor losses (which occur at bend, orifice, change of diameter, etc.) and large losses, like the loss along a pipeline.

In particular, if we apply Bernoulli's equation to our system, figure 1, taking into account that the pressure on the free surface and at exit orifice are the atmospheric pressure  $P_0$ , we have:

$$\frac{1}{2g} u_2^2 + \frac{1}{2g} k u_2^2 + \Delta Z = h + \frac{\alpha_1}{2g} u_1^2, \quad (4)$$

here the term  $k$  is the coefficient of minor losses at the exit orifice, point 2 in our case,<sup>1,2</sup> and  $\Delta Z$  represents the part of the energy loss independent of velocity. As justified in appendix A, the velocity coefficient  $\alpha_2$  is equal to 1 for the exit jet through the orifice. The variable  $h$  ( $=z_1 - z_2$ ) represents the height of the free surface relative to the position of the exit orifice.

The velocities  $u_1$  and  $u_2$  are related by the continuity equation (conservation of mass). For an incompressible fluid, ( $\rho_1 = \rho_2$ ) we have:

$$d_2^2 \cdot C_v \cdot u_2 = d_1^2 \cdot u_1, \quad (5)$$

here  $d_1$  and  $d_2$  represent the diameters of the vessel and the orifice respectively.  $C_v$  is the coefficient of Vena Contracta,<sup>1,2,3</sup> that is related to the fact that the cross section of the exit jet in general is smaller than that of the exit orifice. Combining Eqs. (4) and (5) we obtain:

$$u_2^2 = \frac{2g \cdot (h - \Delta Z)}{\left[ 1 + k - \alpha_1 \cdot \left( \frac{d_2}{d_1} \right)^4 C_v^2 \right]}. \quad (6)$$

Since in our case,  $d_1$  is considerably larger than  $d_2$  ( $d_2/d_1 \approx 0.03$ ) and  $C_v < 1$ ; expression (6) can be simplified to:

$$u_2^2 \cong \frac{2g}{[1+k]} (h - \Delta Z). \quad (7)$$

The implication of this expression is that, if we can measure the velocity of the exit jet  $u_2$  as a function of  $h$ , according to our model, we should expect a linear relationship between  $u_2^2$  and  $h$ . The slope and intersection with the axis would allow us to determine experimentally the coefficients  $k$  and  $\Delta Z$ .

From this same model, it is possible to determine the motion of the free surface ( $u_1$ ) and the evacuation time  $t_e$ .<sup>6</sup> Combining Eqs. (5) and (7) we have:

$$u_1 = -\frac{dh}{dt} = -C_v \left( \frac{d_2}{d_1} \right)^2 \sqrt{2g\mu} \sqrt{h - \Delta Z}, \quad (8)$$

where we have introduced the coefficient  $\mu=1/(1+k)$ . The sign in this equation is related with the orientation chosen to define the positive direction of  $h$  and  $u_1$ . Introducing a new constant  $A$  as:

$$A = C_v \left( \frac{d_2}{d_1} \right)^2 \sqrt{2g\mu} \quad (9)$$

we can write a differential equation for  $h$ , which is:

$$\frac{dh}{\sqrt{h - \Delta Z}} = -A dt. \quad (10)$$

This expression can be easily integrated to give:

$$\frac{\sqrt{h - \Delta Z}}{\sqrt{h_0 - \Delta Z}} = \left( 1 - \frac{t}{t_e} \right), \quad (11)$$

where  $h_0$  is the height of the free surface for  $t=0$  and  $t_e$  is the empty time, which, according to Eqs. (9) and (11), is given by:

$$t_e = \frac{\sqrt{h_0 - \Delta Z}}{C_v \left( \frac{d_2}{d_1} \right)^2 \sqrt{2g\mu}}. \quad (12)$$

Therefore, by measuring  $h$  as a function of time, if the model presented here is correct, the plot of  $\sqrt{h - \Delta Z}$  as a function of  $t$  should be linear. The parameters of this straight line could provide us the value of  $t_e$  and allow us to find the value of  $C_v$ . For  $t \approx t_e$ , there is still some fluid above the orifice ( $h \approx \Delta Z$ ), but there is no jet, the liquid just leaks out of the tank.

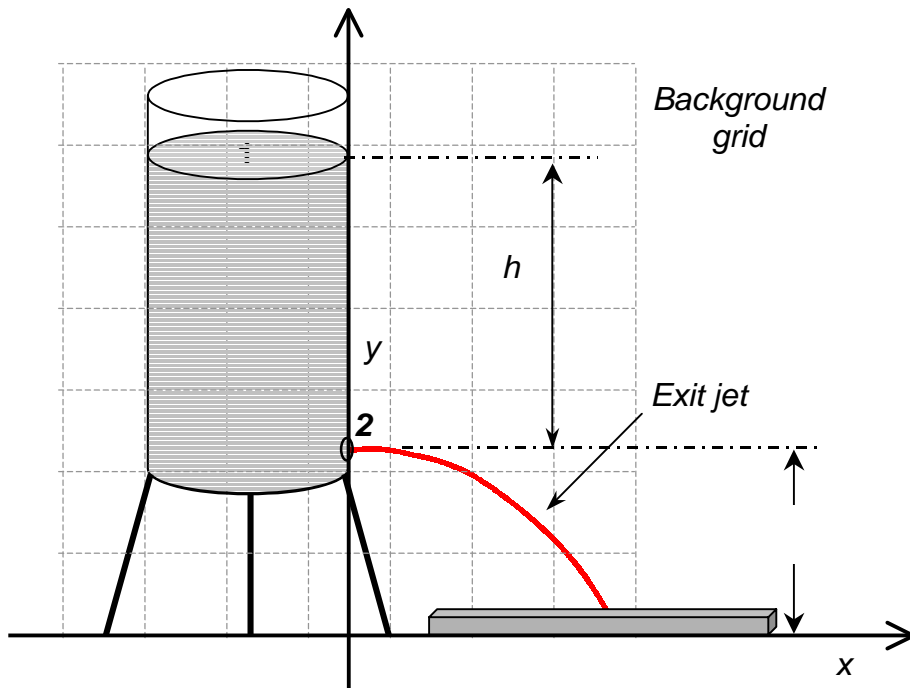
To determine the value of  $u_2$ , we shall assume that the motion of the water particles that comprise the jet, follow the same equation of motion as those of the horizontal projectile with initial velocity  $u_2$ , so we have that:

$$x(t) = u_2 t \quad \text{and} \quad y(t) = H - \frac{1}{2} g t^2, \quad (13)$$

where  $H$  is the height of the exit orifice relative to the origin of coordinate chosen, see figure 2. Combining these two expressions, we obtain the equation for the trajectory of the jet:

$$y(x) = H - \frac{1}{u_2^2} \left( \frac{g}{2} x^2 \right). \quad (14)$$

Therefore, if we can determine experimentally the form of the trajectory of the exit jet of water, by fitting expression (14) to the experimental data, we could obtain the value of  $u_2$ .



**Figure 2:** Schematic diagram of the experimental setup

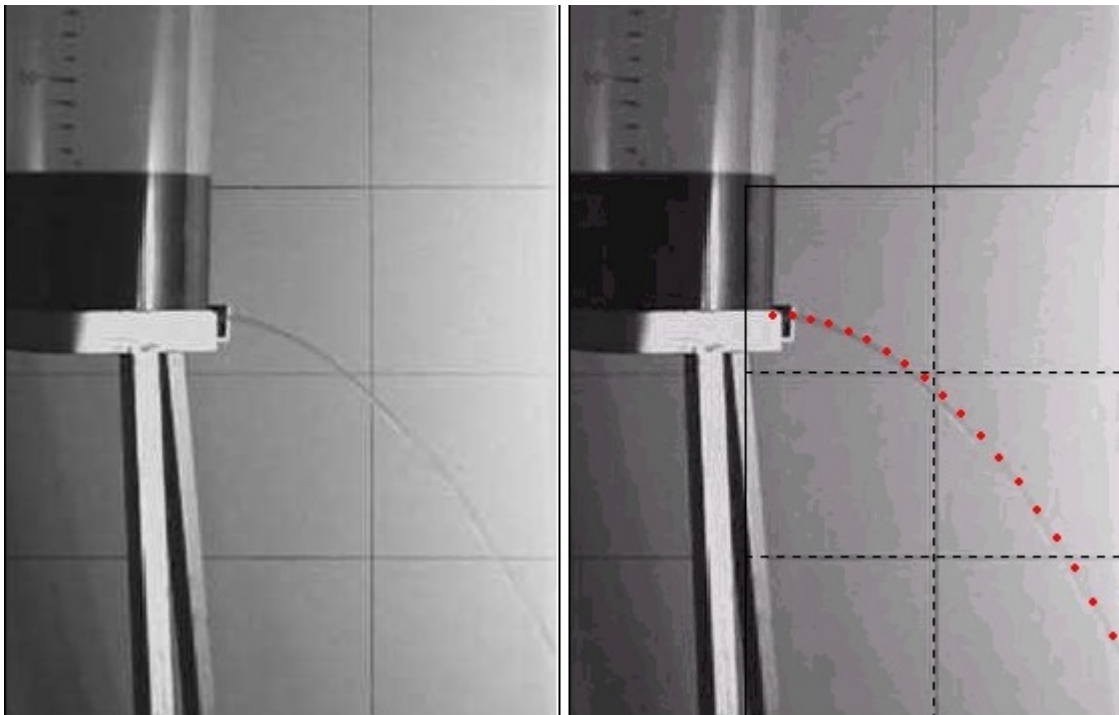
### III. Experiment

The experimental setup consisted of a transparent cylindrical tank with a lateral drainage orifice close to the bottom and a web cam connected to a computer. The vessel was 11 cm in diameter and 25 cm high, the diameter of the orifice was 3 mm and the wall thickness of the vessel was 2 mm. The tank was filled with tap water with a few drops of blue ink to facilitate the visualization of the exit jet. The vessel was positioned in front of a board on which we had drawn a grid with lines every 10 cm to provide an absolute scale for the photographs. The exit jet was parallel to the gridded board. The web cam was placed just in front of the vessel, at about 1.5 m. In this manner we were able to photograph the jet of water with the gridded board in the background. Pictures were taken every time the free surface dropped by about 1 cm. Clearly a digital camera could also be use for this purpose. In this manner, each photograph recorded the height of water  $h$  in the cylinder and the trajectory of the exit jet of water from the orifice. To determine the trajectory of the jet it is possible to use almost any graphics program, such as Microsoft Photo Editor®, to obtain the location of any point in the picture. Furthermore, using the background grid, it is possible to transform the coordinate of any pixel in the photo to the grid coordinate. Another alternative, is to use the picture as a background of a plot, as described bellow.

The procedure we followed to obtain  $u_2$  from the experimental data, consisted in overlapping the digital photograph of the jet to a Microsoft Excel® graph, with a transparent background.<sup>7</sup> The grid of the Excel graph is set to coincide with the mesh grid

used in the background of the picture. The graph is moved and stretched so that the mesh grid of the picture exactly coincides with the corresponding grid in the Excel graph. Once this situation is achieved, the origin (vertex) of the theoretical parabola, described by Equation (14), is chosen to coincide with the exit orifice. Then the value of  $u_2$  is varied so that the theoretical curve described by Eq. (14) coincides with the experimentally observed trajectory. Figure 3 is an example that illustrates this procedure.

We also tracked the position of the free surface of the water in time. To facilitate the measurement of the height of liquid  $h$ , we drew horizontal marks every 0.5 cm, starting at the position of the drainage orifice. We used a stop watch to measure the time it took for the free surface of the water to reach each horizontal mark. In this manner we measured  $h$  as a function of time.

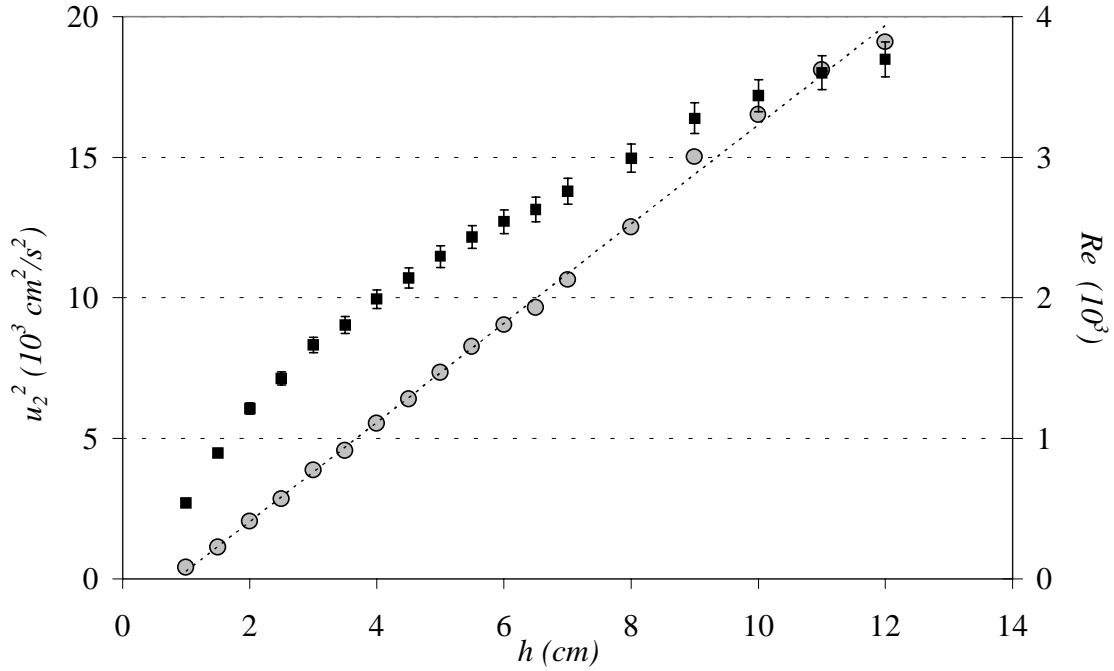


**Figure 3:** a) Actual digital photograph of the exit jet. b) The same photograph in the background overlapped with a plot of the theoretical trajectory (dotted line) of the jet as described by expression (14).

#### IV. Results and Discussion

In figure 3.a) we present an actual digital photograph of the vessel and the exit jet. In figure 3.b) we show the same digital photograph of the exit jet in the background overlapped with a plot of the theoretical trajectory of the jet as described by expression (14). By adjusting the value of  $u_2$ , we “fitted” the theoretical trajectory to the actual trajectory of the jet. The grid lines in the Excel plot were chosen to coincide with the grid lines in the background of the picture<sup>7</sup>. This figure also shows, that the trajectory of the jet is very well reproduced by Eq. (14). This agreement is a clear indication that the

liquid particles that conform the jet follow the same trajectory as the solid particles do. Therefore the liquid elements of volume are described by the same physical laws of mechanics that govern the motion of solids. This result may be useful to confronting the Aristotelian misconception, still prevalent in some students, that liquids and solids follow different laws.

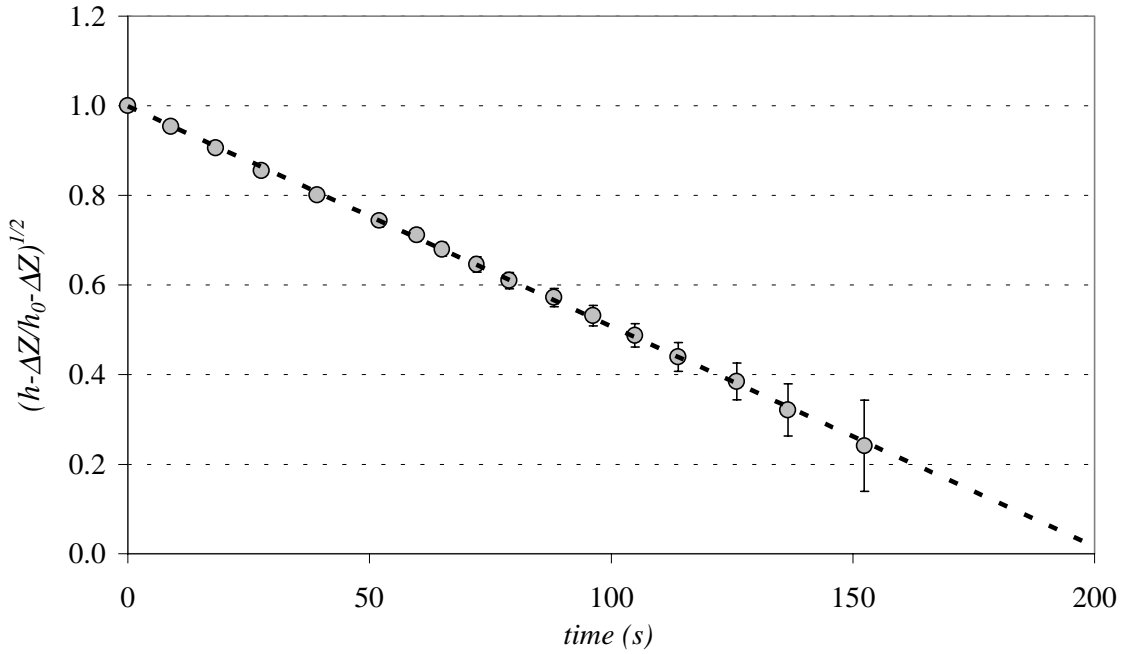


**Figure 4:** The circles represent the experimental results of  $u_2^2$  as a function of the height  $h$ , the straight line is a linear fit to the data. The square symbols indicate the values of the Reynolds number (right vertical scale) at the orifice for the different values of  $h$  used in our experiment. By choosing a taller cylinder it is possible to further explore turbulent regime.

In figure 4, we present the plot of  $u_2^2$  as a function of  $h$ . The fact that the relationship between  $u_2^2$  and  $h$  is linear, indicates that the main assumptions of the model proposed, expressed by Eqs. (6) and (7), are in good agreement with the experimental results. Furthermore, by fitting the theoretical expectation, Eq. (7), to the data, we can obtain the values of the parameters  $\Delta z$  and  $k$ . In table 1 we present the results of these parameters for different runs of the experiment. We performed several runs of the same experiment to test the consistency of our results and the robustness of the parameters of the model.

In figure 5 we present the results for the height of the free surface  $h$  as a function of time. To test the validity of our model, expression (11), we plotted the modified variable  $\sqrt{(h - \Delta Z)/(h_0 - \Delta Z)}$  as a function of time. The fact that this plot has a clear linear trend, indicates the validity of the model to describe the physics of the problem. Furthermore, the parameter of the fitted line allows us to obtain the values of the empty time  $t_e$  and the coefficient of vena contracta  $C_v$ . In table 1 we present the values of the relevant parameters of the model for several runs of the experiment. We notice that the

values of the measured parameters are consistent for the different runs of the experiment. Furthermore, the value of  $k$  (coefficient of minor losses) and  $C_v$  are consistent with the values reported in the literature.<sup>1,2,3</sup> In figure 4, on the right vertical axis we plotted the value of the Reynolds number,  $Re$ , at the orifice for the different values of  $h$  used in our study. We see that in our experiment the value of  $Re$  varies between 4000 and nearly 0, indicating that the flow regime spans beginning of the turbulent, transitional and laminar flows. The fact that the model is able to reproduce the experimental data in all these regions of  $Re$ , indicates that hypothesis made in our model, namely the type of losses proposed and the expanded Bernoulli's equation, are valid for these regimes of flows.



**Figure 5:** The circles represent the experimental results of  $\sqrt{(h - \Delta Z)/(h_0 - \Delta Z)}$  as a function of time, the straight line is a fit to the data.

Parameter	Run 1	Run 2	Run 3
$\Delta Z$ (cm)	$0.9 \pm 0.2$	$0.9 \pm 0.2$	$0.9 \pm 0.2$
$\mu$	$0.91 \pm 0.01$	$0.90 \pm 0.01$	$0.90 \pm 0.01$
$k$	$0.09 \pm 0.01$	$0.11 \pm 0.01$	$0.11 \pm 0.01$
$C_v$	$0.46 \pm 0.03$	$0.47 \pm 0.03$	$0.44 \pm 0.03$

Table 1: Parameters obtained for several runs. The values of  $\mu$  and  $C_v$  are consistent with those reported in the literature.<sup>1,2,3</sup>



## V. Conclusions

The experiment carried out in this study is simple and inexpensive. Also it is accessible to beginner students and it clearly illustrates the importance and usefulness of Bernoulli's equation for real fluids including energy losses, for a wide range of Reynolds numbers. The validity of the type of energy losses proposed and the expanded Bernoulli's equation, spans the beginning of the turbulent regime, as well as the transitional and laminar regimes. All that is needed to further explore the turbulent regime is a taller cylinder.

The experiment also verifies that liquid particles follow the same trajectory as solid particles, indicating that they obey the same physical laws.

The proposed model, based on the extended Bernoulli's equation, is adequate to describe qualitatively and quantitatively the physics of drainage of a vessel. Furthermore, the fitting of the theoretical model to the data allows us to extract the relevant parameters of the model, namely the coefficients  $\Delta Z$  and  $k$ , of Eq. (4), and the coefficient of vena contracta  $C_v$ . Our results are consistent with values of these parameters that are independent of the Reynolds number in the region explored in this study.

The motion of the free surface is well reproduced by the model, in particular its dependence with time. Furthermore, this study allows us to measure the coefficient of vena contracta in the exit jet and the coefficient of losses in our system.

## VI. Acknowledgment

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## Appendix A

When a laminar flow enters to a pipe, it does not immediately develop the parabolic velocity profile that prevails well inside the pipe. There is a transition region or entrance length<sup>1,3</sup> ( $L_e$ ) over which the velocity profile changes from a rather planar to the fully developed parabolic profile. There is a semi empirical relation that allows us to estimate this entrance length given by:<sup>8</sup>

$$L_e = C_e \cdot Re \cdot d, \quad (A1)$$

here  $Re(=\rho \cdot u \cdot d / \eta)$  is the Reynolds number,  $d$  the inner diameter of the pipe and  $C_e$  is a constant. Several authors have proposed different values for this constant,<sup>1,3,8</sup> but all range between 0.029 and 0.06 for laminar flows. If we apply this expression to the exit orifice, in all cases studied here, we obtained values of  $L_e$  that were much larger than the wall thickness. Therefore, the velocity profile of the exit jet can be regarded as planar in

this case, i.e.  $\alpha_2 \approx 1$ . On the other hand, for the case of the cylinder, the average velocity is so small, ( $u_1 \approx 10^{-2}$  cm/s) the value of  $\alpha_1$  is approximately 2.

### Figure Captions

**Figure 1:** Schematic diagram of the experimental setup.

**Figure 2:** Schematic diagram of the experimental setup

**Figure 3:** a) Actual digital photograph of the exit jet. b) The same photograph in the background overlapped with a plot of the theoretical trajectory (dotted line) of the jet as described by expression (14).

**Figure 4:** The circles represent the experimental results of  $u_2^2$  as a function of the height  $h$ , the straight line is a linear fit to the data. The square symbols indicate the values of the Reynolds number (right vertical scale) at the orifice for the different values of  $h$  used in our experiment. By choosing a taller cylinder it is possible to further explore turbulent regime.

**Figure 5:** The circles represent the experimental results of  $\sqrt{(h - \Delta Z)/(h_0 - \Delta Z)}$  as a function of time; the straight line is a fit to the data.

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<sup>2</sup> V. L. Streeter and E. D. Wylie, *Fluid Mechanics*, 8th edition (Mc Graw-Hill, New York, 1985), Chap. 8.

<sup>3</sup> B. Nekrasov, *Hidráulica* (in Spanish), (3rd edition, MIR, Moscow, 1968), Chap. 6 and 9.

<sup>4</sup> C. E. Synolakis and H. S. Badeer, "On combining the Bernoulli and Poiseuille equation-A plea to authors of college physics texts", *Am. J. Phys.* 57, 1013 - 1019 (1989)

<sup>5</sup> C. Waltham, S. Bendall and A. Kotlicki, "Bernoulli levitation", *Am. J. Phys.* 71, 176 - 179 (2003)

<sup>6</sup> J. N. Libii, "Mechanics of the slow draining of a large tank under gravity", *Am. J. Phys.* 71, 1204 - 1207 (2003).

<sup>7</sup> Examples of the Excel files illustrating this procedure can be downloaded from [www.fisicareCreativa.com](http://www.fisicareCreativa.com). Web site dedicated to experimental physics. It contains experimental projects using new technologies and reports of experiments performed by undergraduate students from different universities of Latin America.

<sup>8</sup> S. G. Kandlikar and L. A. Campbell, "Effect of entrance condition on frictional losses and transition to turbulence", Proceedings of IMECE2002, ASME International Mechanical Engineering Congress & Exposition, November 17-22, 2002, Louisiana, New Orleans.