

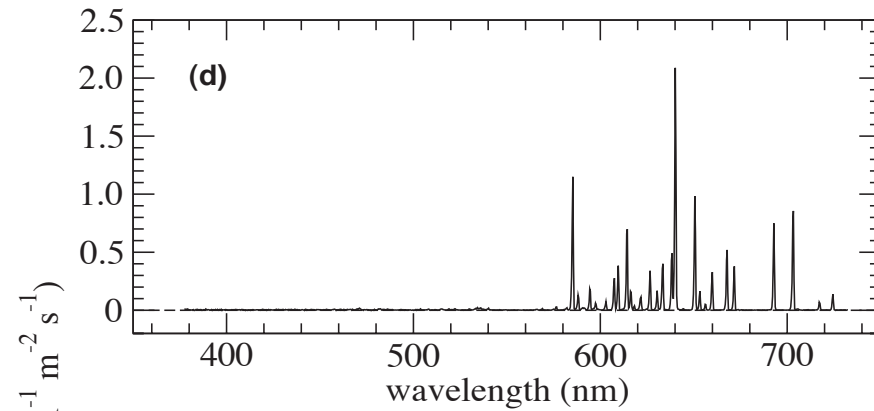
# Introduction to the plasma spectroscopy

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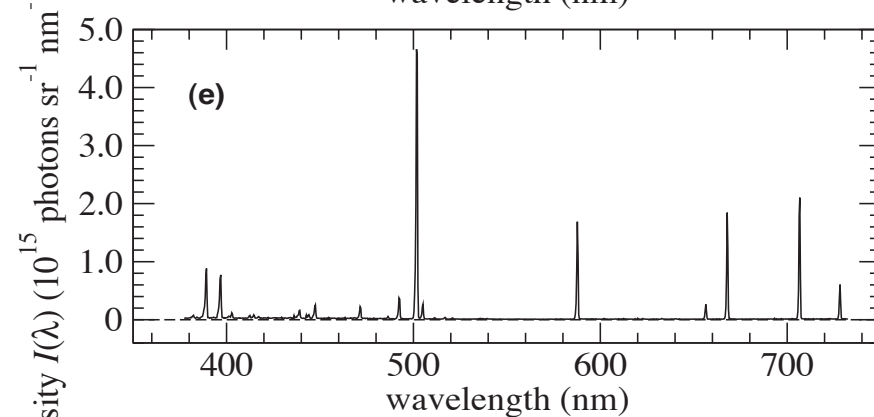
- Einstein coefficients
- energy state designation under  $L$ - $S$  coupling scheme
- population mechanism for excited states — collisional-radiative model
- effect of external field — perturbation theory

**introduction**

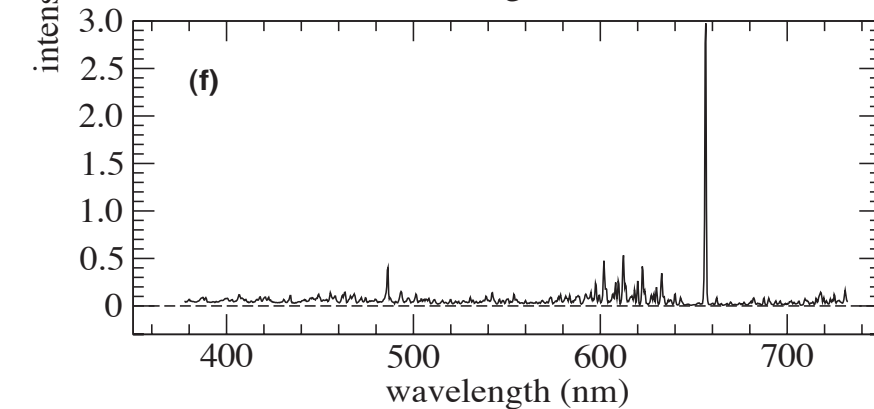
# glow discharge in LHD



Ne

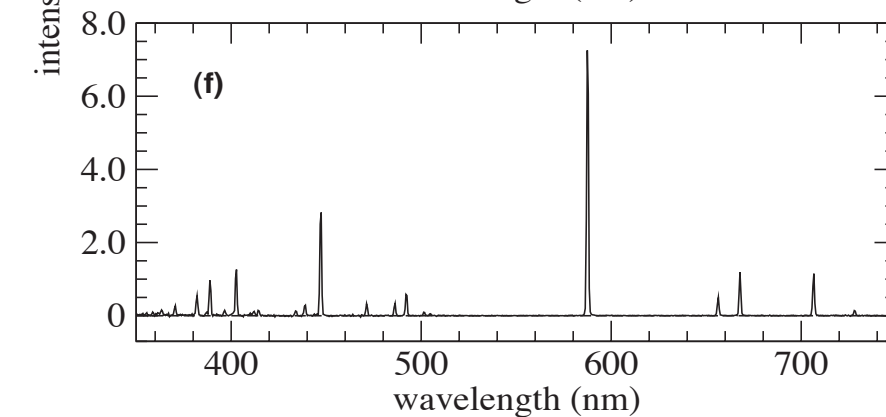
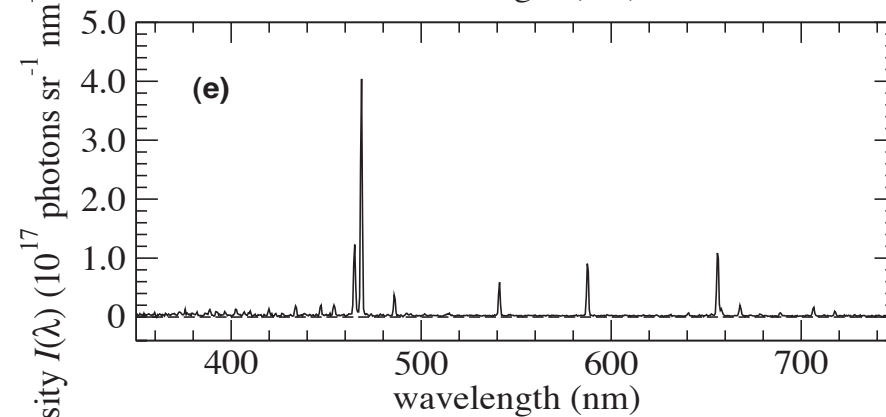
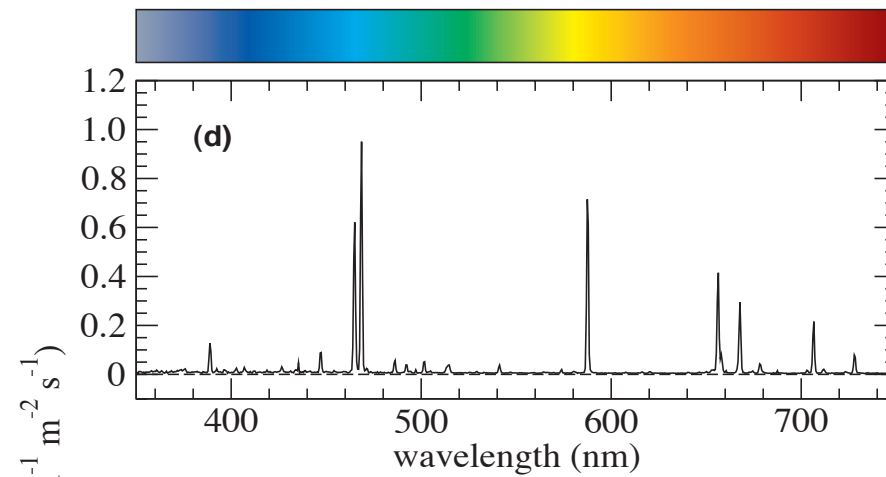
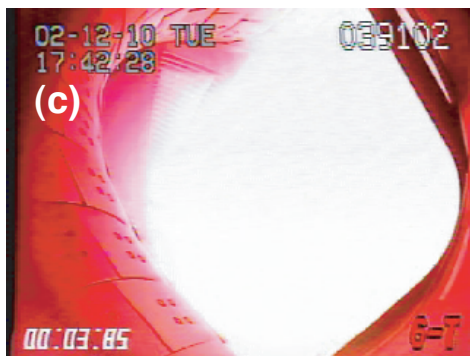
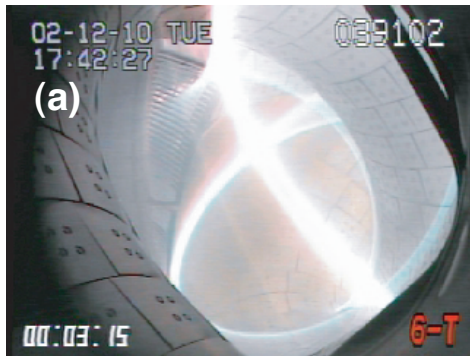


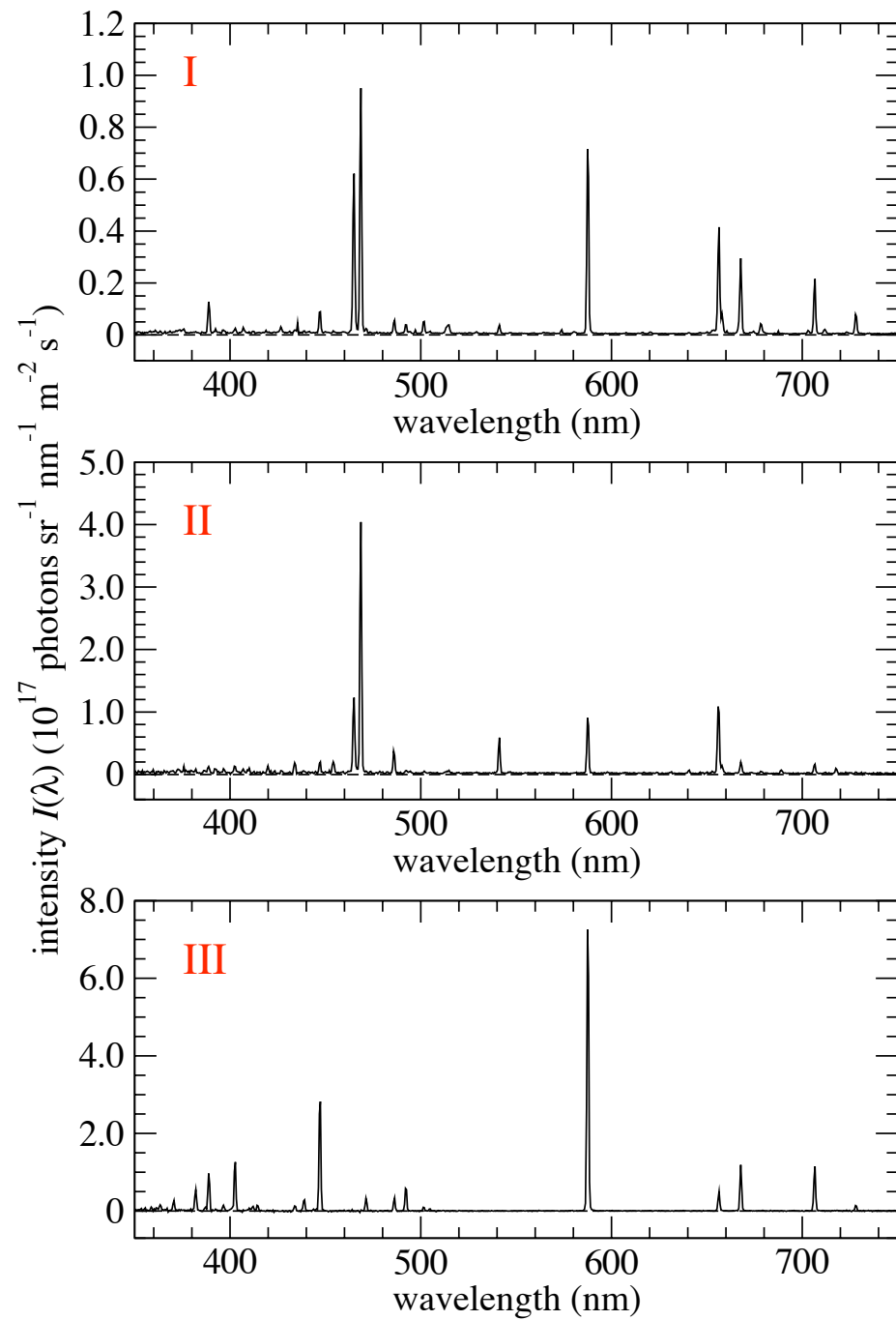
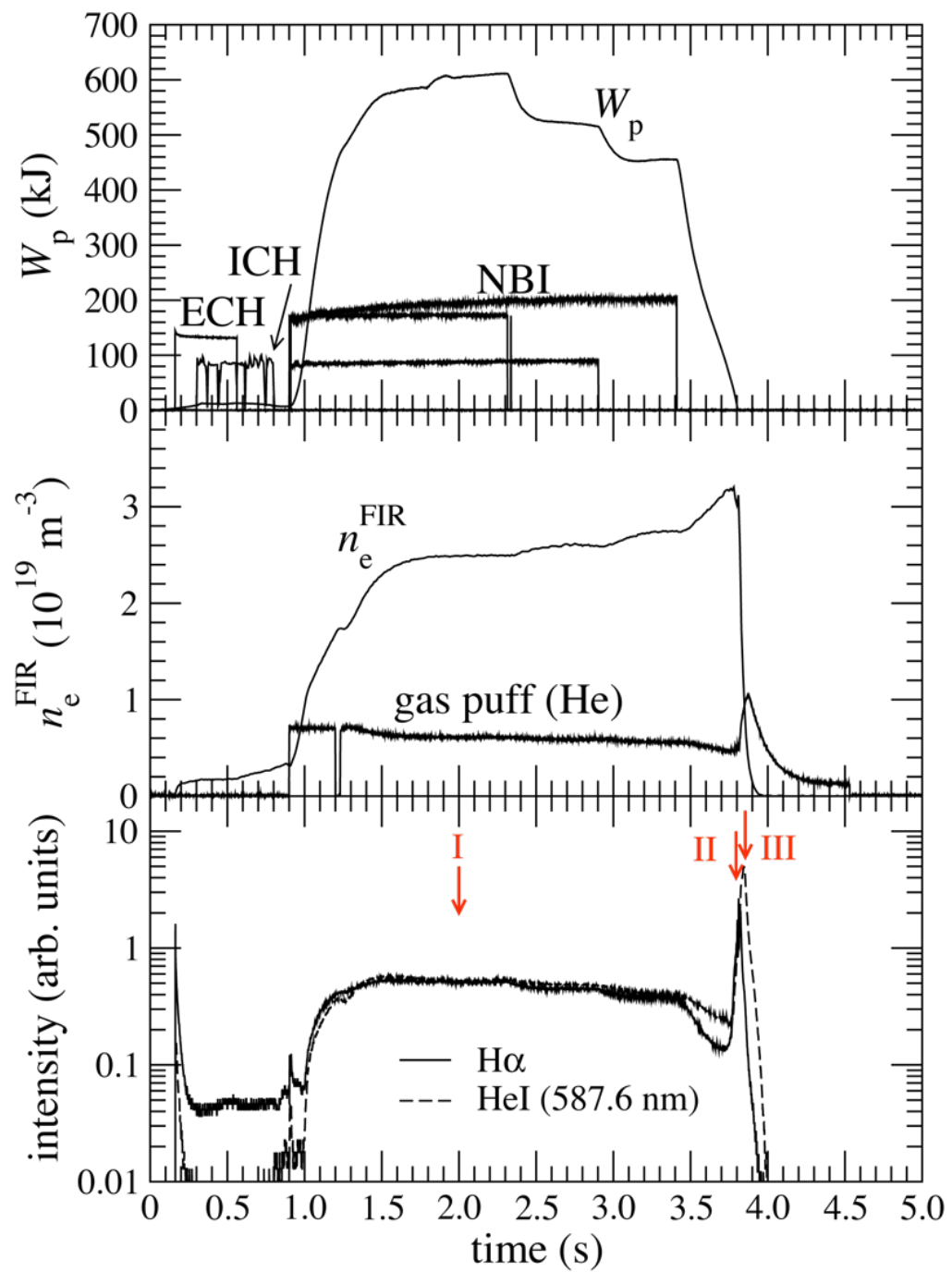
He



H<sub>2</sub>

# main discharge with helium gas



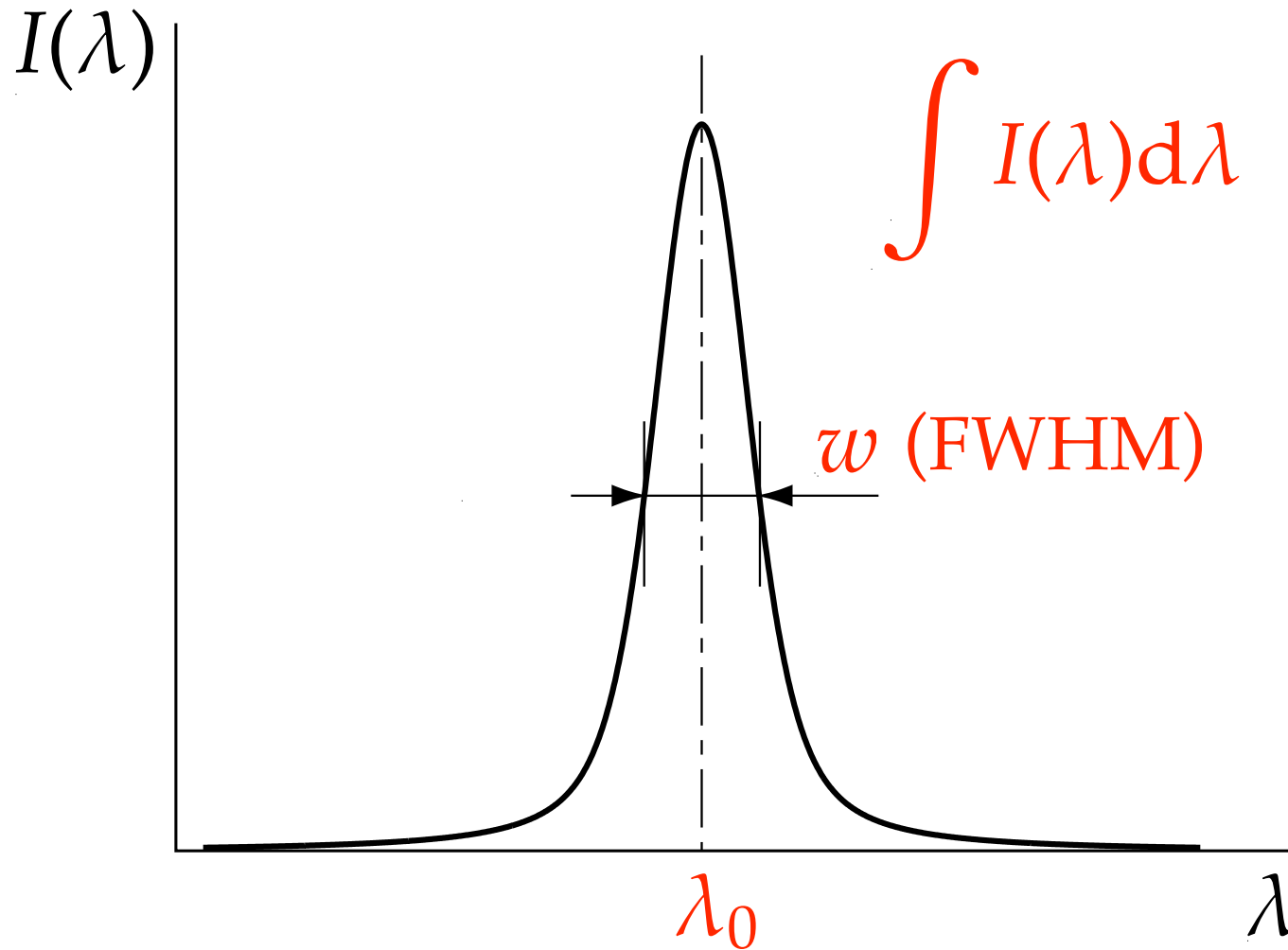


- plasma shows various **colors** with different working gas
- quantitative treatment of **color** is, however, difficult
- **spectrum** is used for quantitative analysis instead

**emission lines**



# properties of emission line



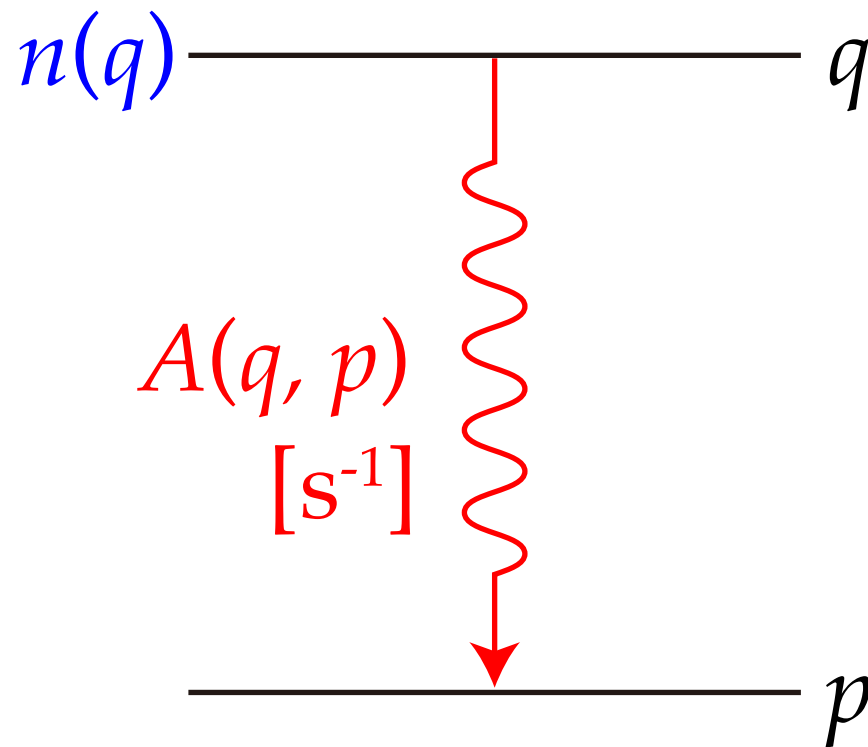
# what can be known?

observable		obtainable
shift		ion velocity
broadening	Doppler	$T_i$
	Stark	$n_e$
splitting	Zeeman	magnetic field
	Stark	electric field
intensity ratios intensity distribution		$T_e, n_e$ ionizing or recombining
intensity		$n_i$

} high resolution measurement  
} low resolution measurement

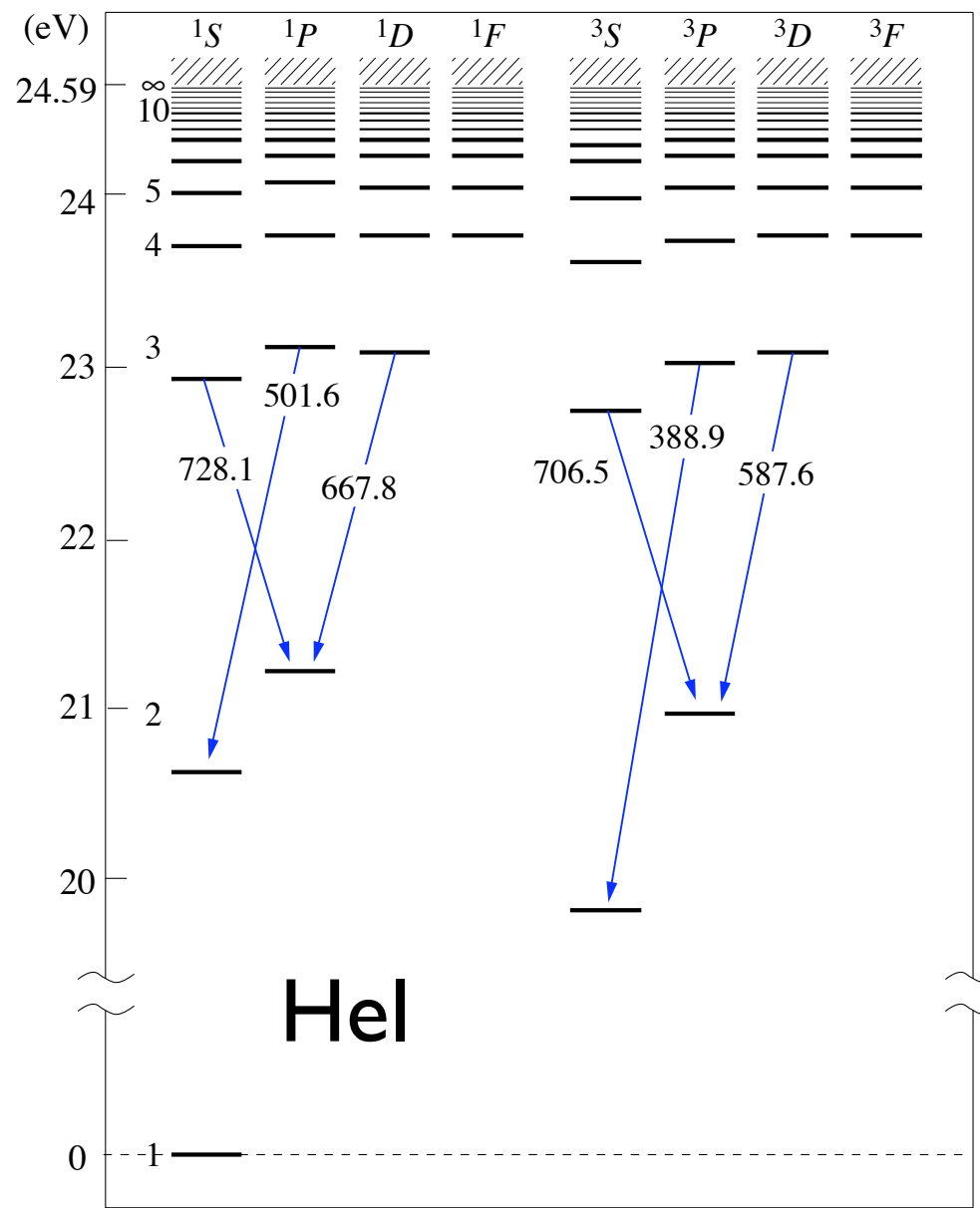
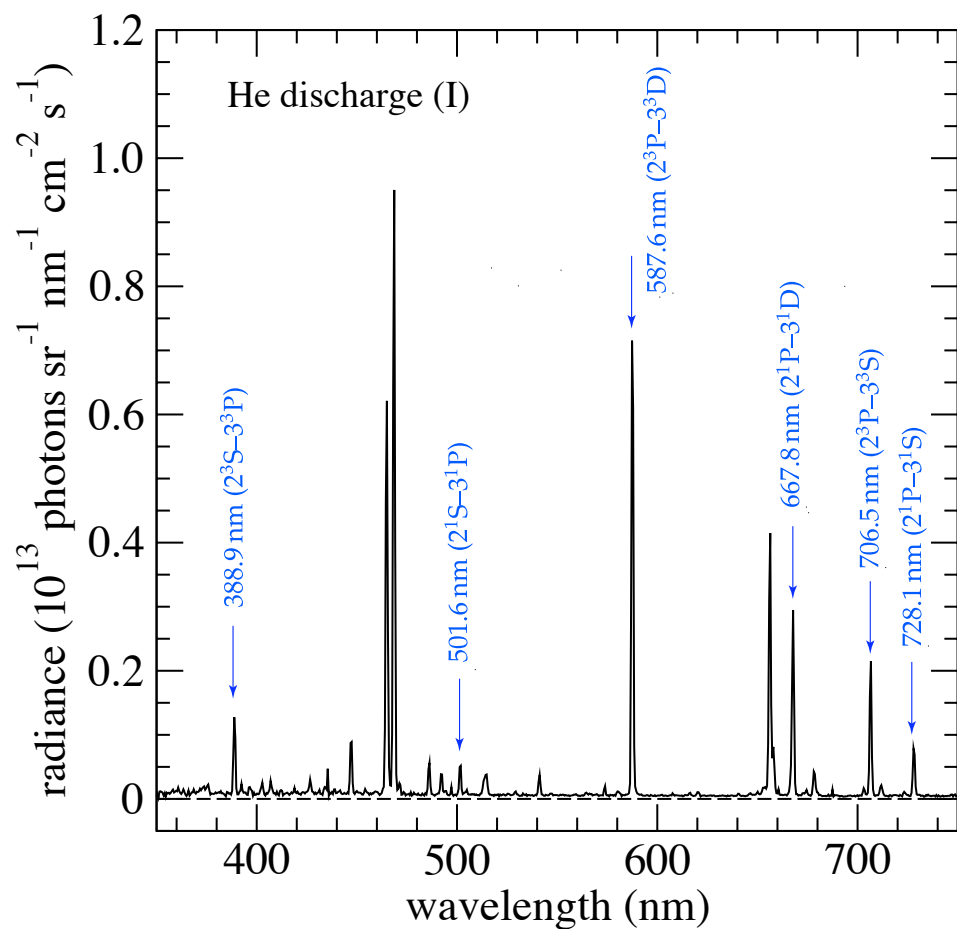
# line intensity

- intensity is the product of **population density** and **spontaneous transition probability** or **Einstein A coefficient**

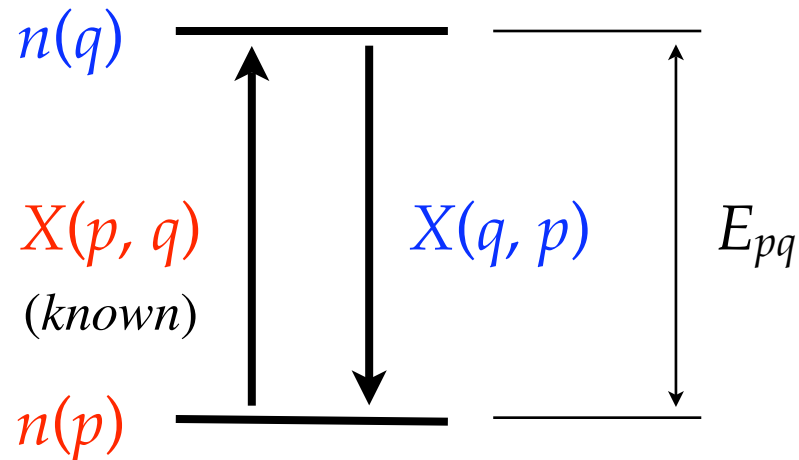


$$I_{qp} = n(q)A(q, p)$$
$$[m^{-3}s^{-1}]$$

A coefficient is constant and proper to individual transitions



# detailed balance



under thermodynamic equilibrium

$$n(p)X(p, q) = n(q)X(q, p)$$

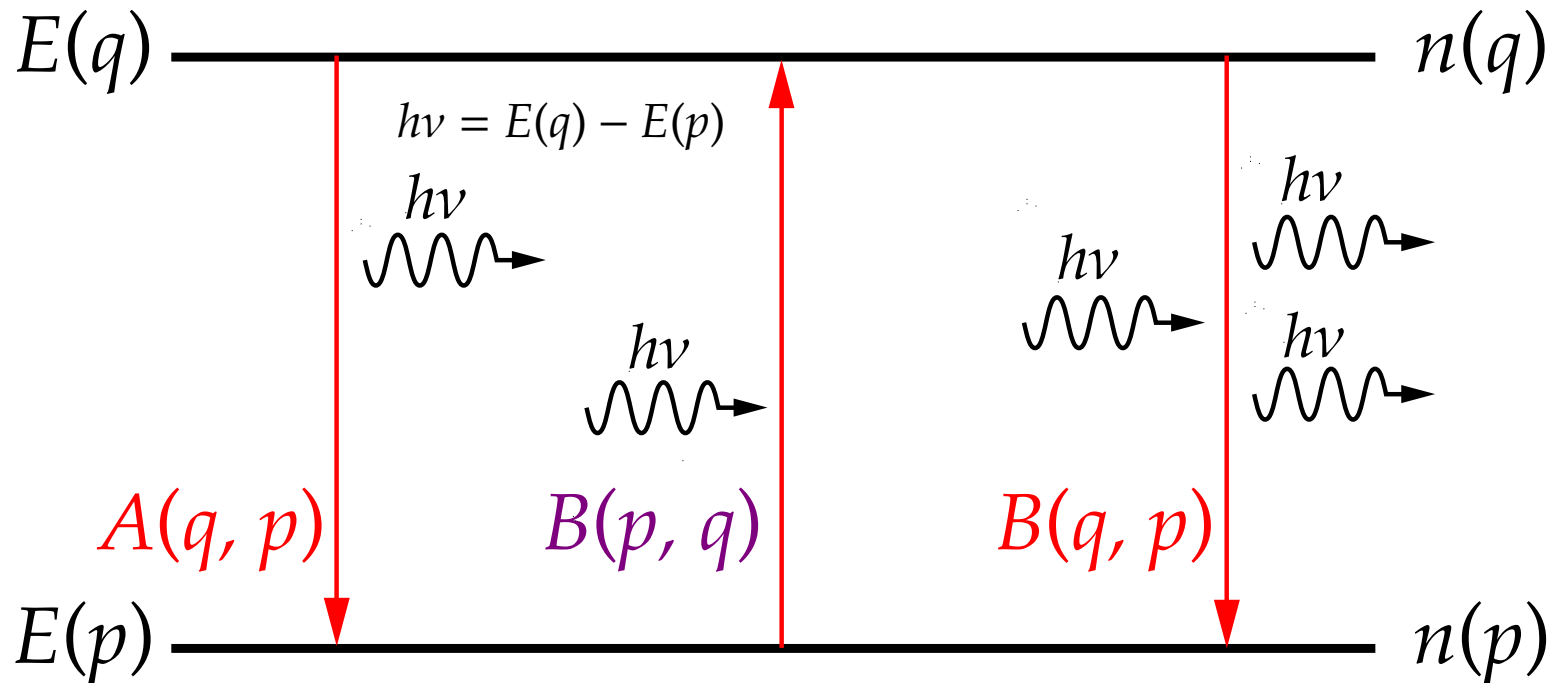
$n(q)$  and  $n(p)$  should obey Boltzmann distribution

$$\frac{n(p)}{n(q)} = \frac{g(p)}{g(q)} \exp\left[\frac{E_{pq}}{kT}\right]$$

$X(q, p)$  must be

$$X(q, p) = X(p, q) \frac{g(p)}{g(q)} \exp\left[\frac{E_{qp}}{kT}\right]$$

# Einstein coefficients



$B(p, q)$  can be calculated, then how  $A(q, p)$  and  $B(q, p)$  are derived?

$$A(q, p)n(q) + B(q, p)n(q)I_\nu = B(p, q)n(p)I_\nu$$

under thermodynamic equilibrium

population ratio is subject to Boltzmann distribution

$$\frac{n(q)}{n(p)} = \frac{g(q)}{g(p)} \exp \left[ -\frac{h\nu}{kT} \right]$$

the balance equation is rewritten as

$$I_\nu = \frac{\frac{A(q,p)}{B(q,p)}}{\frac{B(p,q)}{B(q,p)} \frac{g(p)}{g(q)} \exp \left( \frac{h\nu}{kT} \right) - 1}$$

this should be equivalent to Planck's black-body equation

$$I_\nu = \frac{2h\nu^3}{c^2} \frac{1}{\exp(h\nu/kT) - 1}$$

finally, interrelationships among **Einstein coefficients** are derived as

$$B(q, p) = \frac{g(p)}{g(q)} B(p, q)$$
$$A(q, p) = \frac{2h\nu^3}{c^2} B(q, p) = \frac{2h\nu^3}{c^2} \frac{g(p)}{g(q)} B(p, q)$$



# level notations

(angular momentum theory)

# electron states for hydrogen

orbital angular momentum quantum number

<i>n</i> \ <i>l</i>	0	1	2	3	4	5
	s (2)	p (6)	d (10)	f (14)	g (18)	h (22)
1 (2)	1s					
2 (8)	2s	2p				
3 (18)	3s	3p	3d			
4 (32)	4s	4p	4d	4f		
5 (50)	5s	5p	5d	5f	5g	
6 (72)	6s	6p	6d	6f	6g	6h

$$m_l = -l, \dots, l$$

$$m_s = \pm 1/2$$

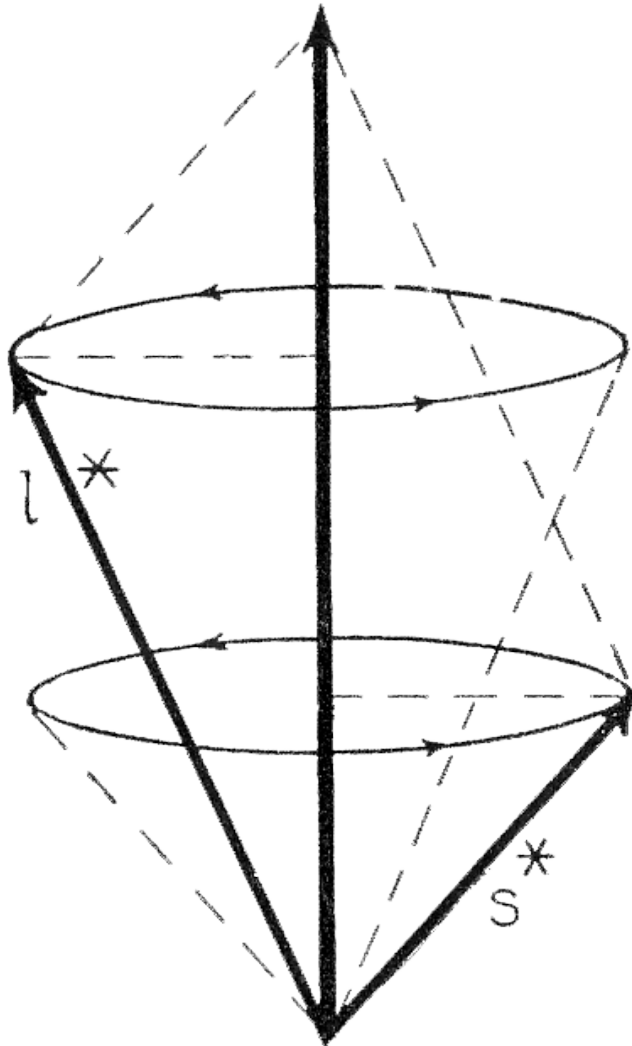
spin quantum number  $s = 1/2$

( ) indicates statistical weight



# coupling of angular momentum

$$j^* \quad j = l + s \quad (|s| = 1/2)$$

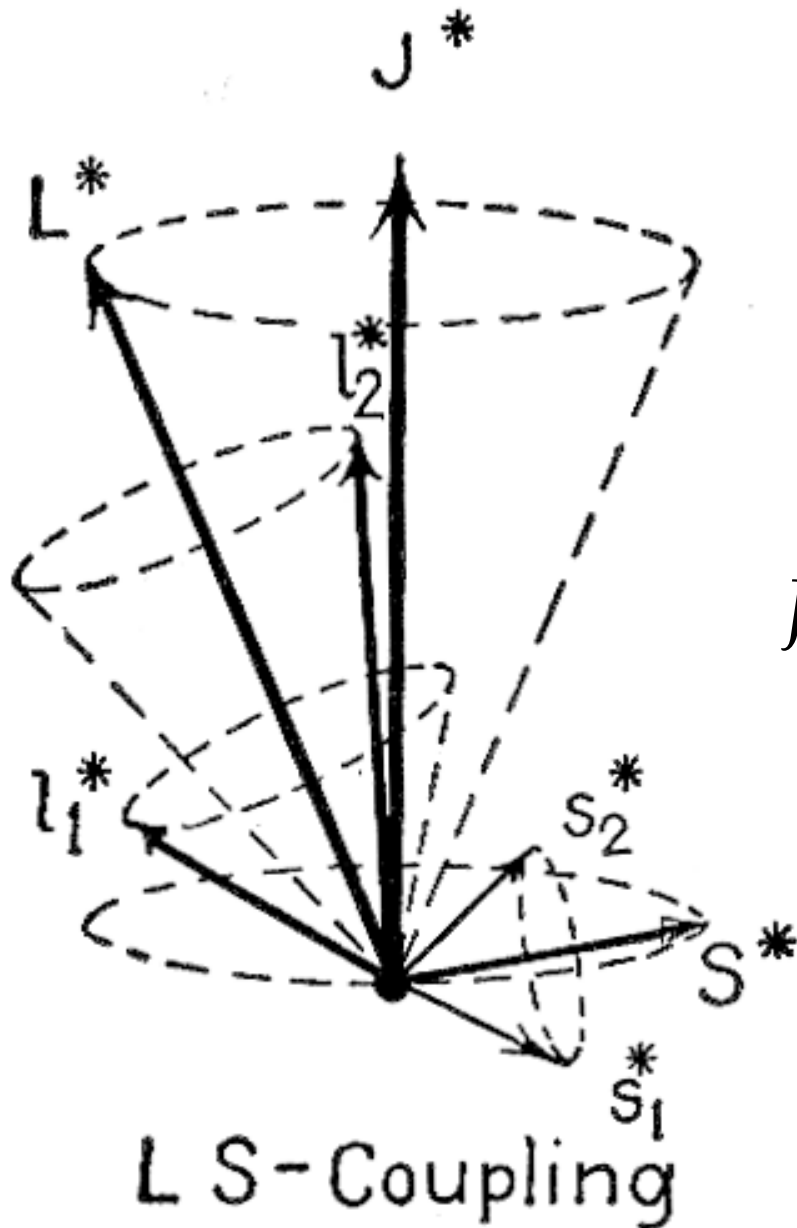


total angular momentum  
quantum number

$$j = \begin{cases} l \pm \frac{1}{2} & \text{for } l \geq 1 \\ \frac{1}{2} & \text{for } l = 0 \end{cases}$$

each  $j$  state has  $2j+1$  magnetic  
substates  $m$  ( $= -j, -j+1, \dots, j$ )

# L-S coupling



$$L = l_1 + l_2$$

$$S = s_1 + s_2$$

$$J = L + S$$

$$J = \begin{cases} L - S, \dots, L + S & \text{for } L \geq S \\ & (2S + 1 \text{ states}) \\ S - L, \dots, S + L & \text{for } L < S \\ & (2L + 1 \text{ states}) \end{cases}$$

each  $J$  state has  $2J+1$  magnetic substates  $M$  ( $= -J, -J+1, \dots, J$ )

# quantum numbers

variables	meanings	values	
$n$	principal	0, 1, 2, ...	
$s$	spin	1/2	
$S$		0, 1/2, 1, 3/2, 2, ...	
$l$	orbital	0, 1, 2, ...	s, p, d, f, ...
$L$			S, P, D, F, ...
$j$	total	0, 1/2, 1, 3/2, 2, ...	
$J$			
$m$	magnetic	0, $\pm 1/2$ , $\pm 1$ , $\pm 3/2$ , $\pm 2$ , ...	
$M$			

# term designation

$$nl^k n'l'^{k'} \dots 2S+1L_J$$

configuration	possible terms	isoelectronic sequence
3d	${}^2D_{3/2,5/2}$	H-like
1s4f	${}^1F_3$ ${}^3F_{2,3,4}$	He-like
$1s^2 2p$	${}^2P_{1/2,3/2}$	Li-like
$1s^2 2s^2$	${}^1S_0$	Be-like
$1s^2 2s 2p({}^3P^o) 3p$	${}^2S_{1/2}$ ${}^2P_{1/2,3/2}$ ${}^2D_{3/2,5/2}$ ${}^4S_{3/2}$ ${}^4P_{1/2,3/2,5/2}$ ${}^4D_{1/2,3/2,5/2,7/2}$	B-like

one electronic configuration could have several terms

# selection rules for transition

$$2S+1L_J - 2S'+1L'_J$$

$$\Delta S = 0$$

$$\Delta L = 0, \pm 1 \text{ (} 0 \rightarrow 0 \text{ excluded)}$$

$$\Delta J = 0, \pm 1 \text{ (} 0 \rightarrow 0 \text{ excluded)}$$

besides the rules for total quantum numbers,  
 $\Delta l = \pm 1$  is always required



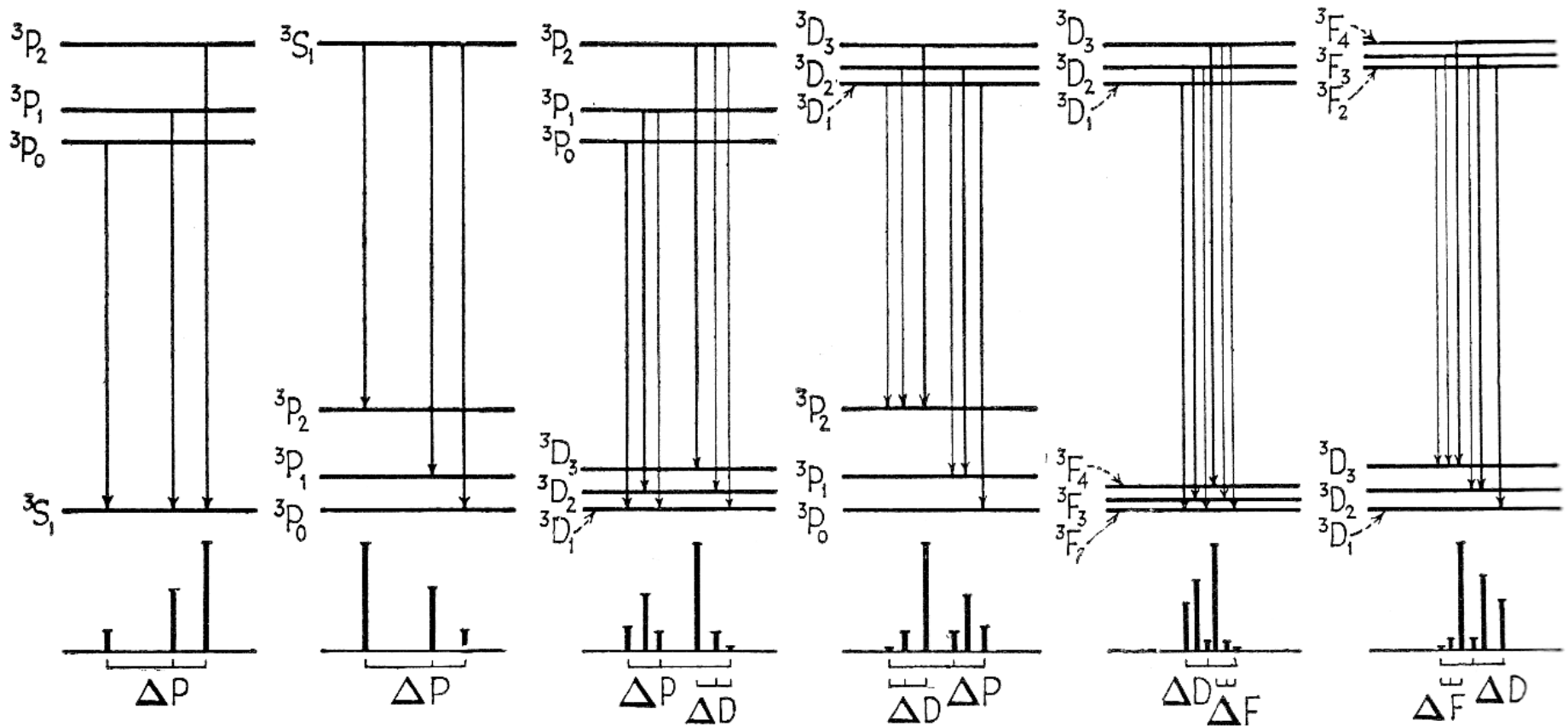
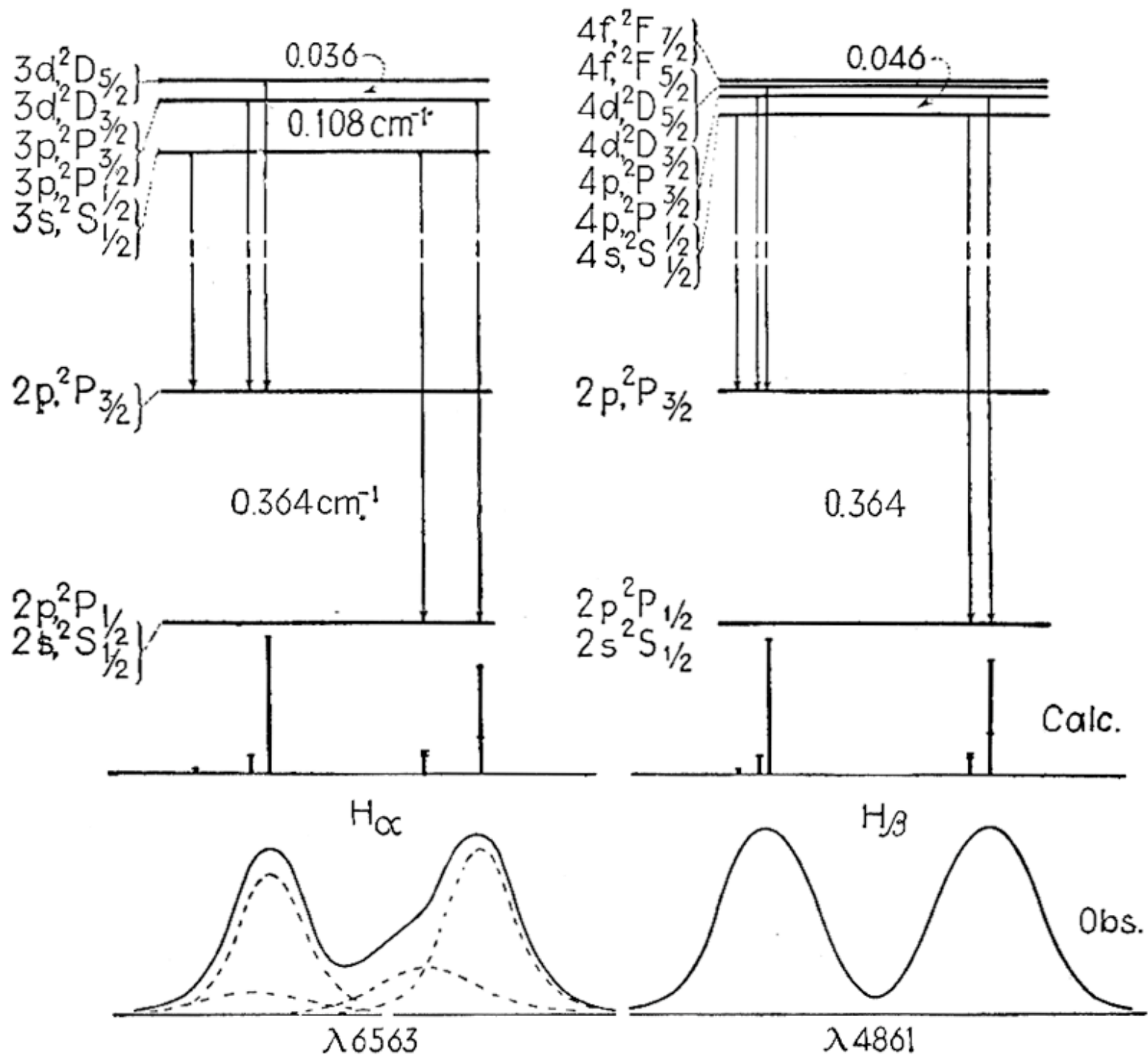


FIG. 11.4.—Triplet-triplet transitions showing selection rules and relative intensities.

(H. E. White, *Introduction to Atomic Spectra*)

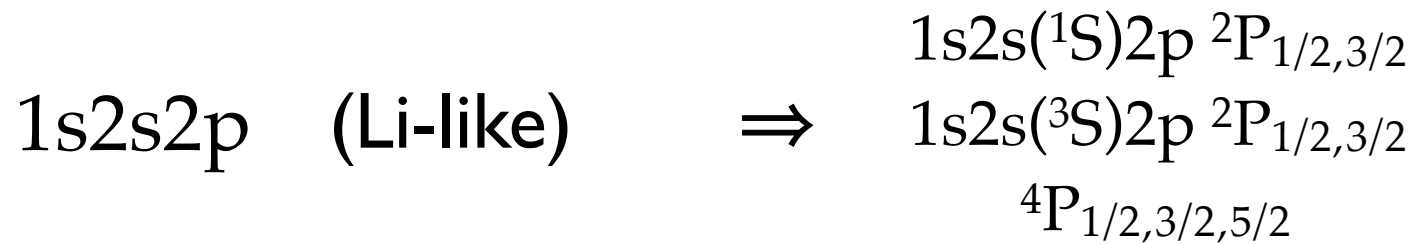
$$I \propto \langle J_i || D || J_k \rangle = (-1)^{S+1+L_i+J_k} \sqrt{g_i g_k} \begin{Bmatrix} L_i & J_i & S \\ J_k & L_k & 1 \end{Bmatrix} \langle L_i || D || L_k \rangle$$



(H. E. White, *Introduction to Atomic Spectra*)

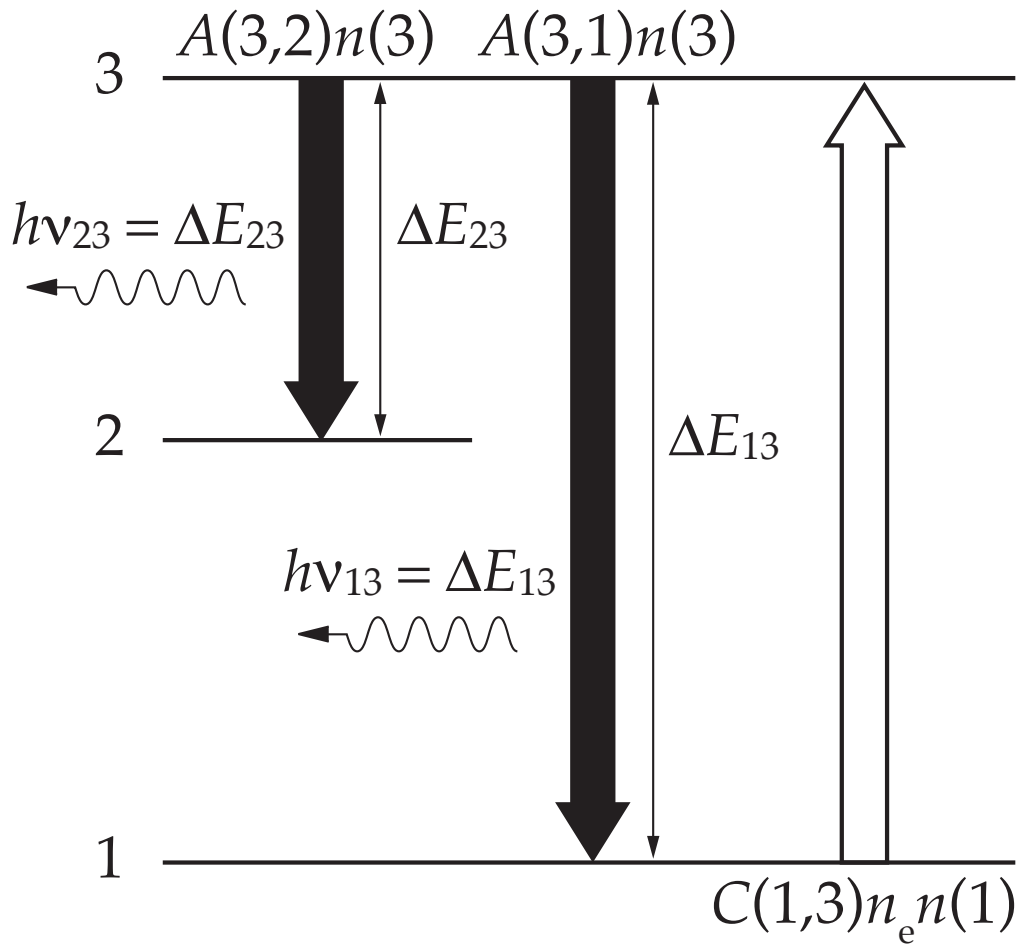
configuration

term(s)



**population mechanisms**

# corona equilibrium



$$C(1,3)n_e n(1) = [A(3,1) + A(3,2)] n(3)$$

more generally

$$C(1,p)n_e n(1) = \sum_{q < p} A(p,q) n(p)$$

$$n(p) = \frac{C(1,p)n_e}{\sum_{q < p} A(p,q)} n(1)$$

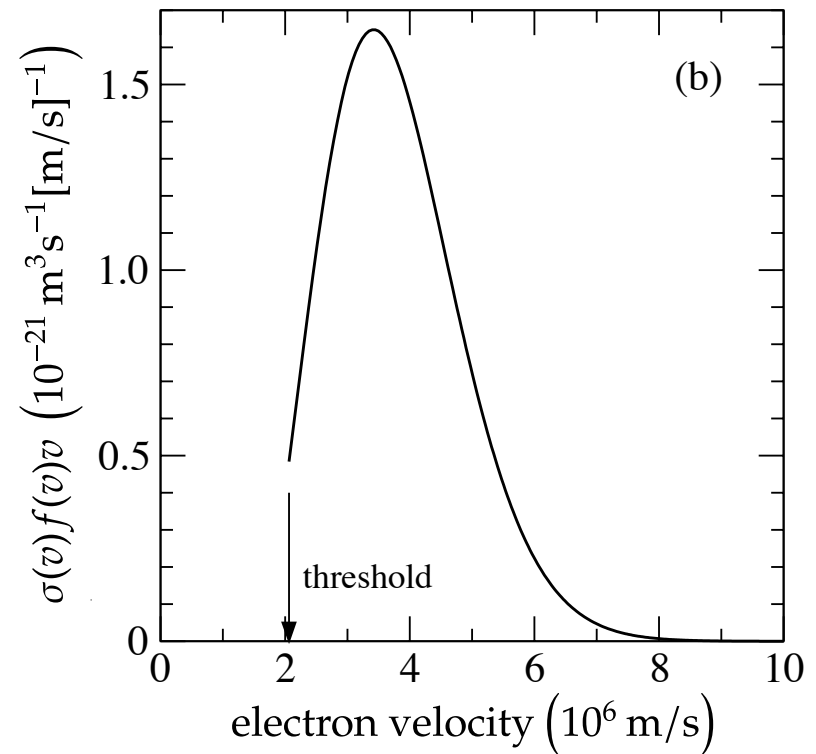
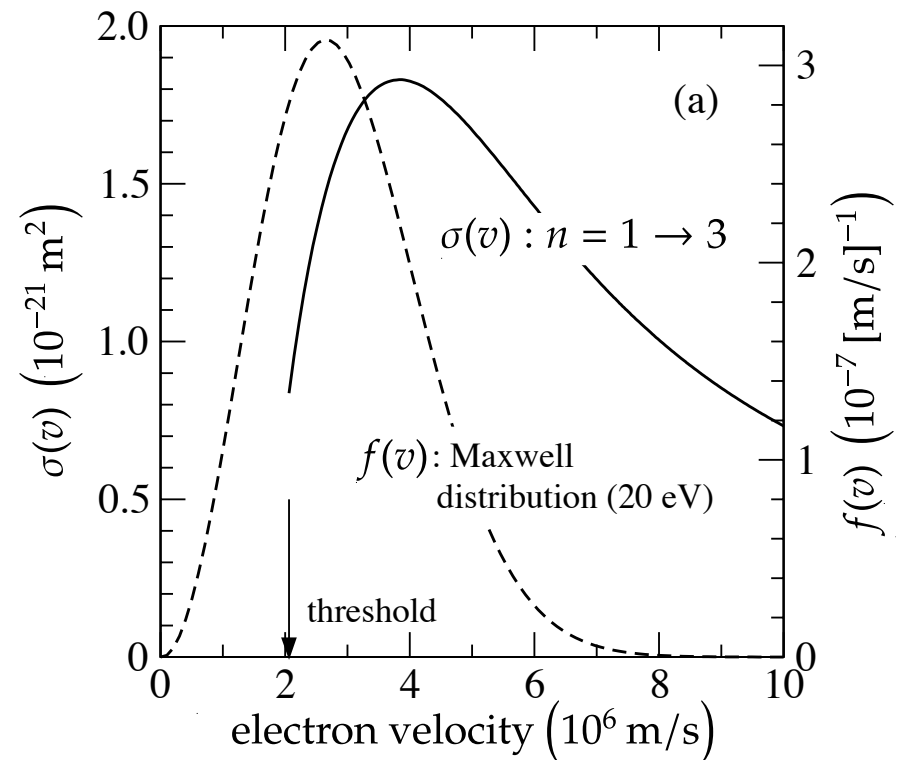
# collisional excitation

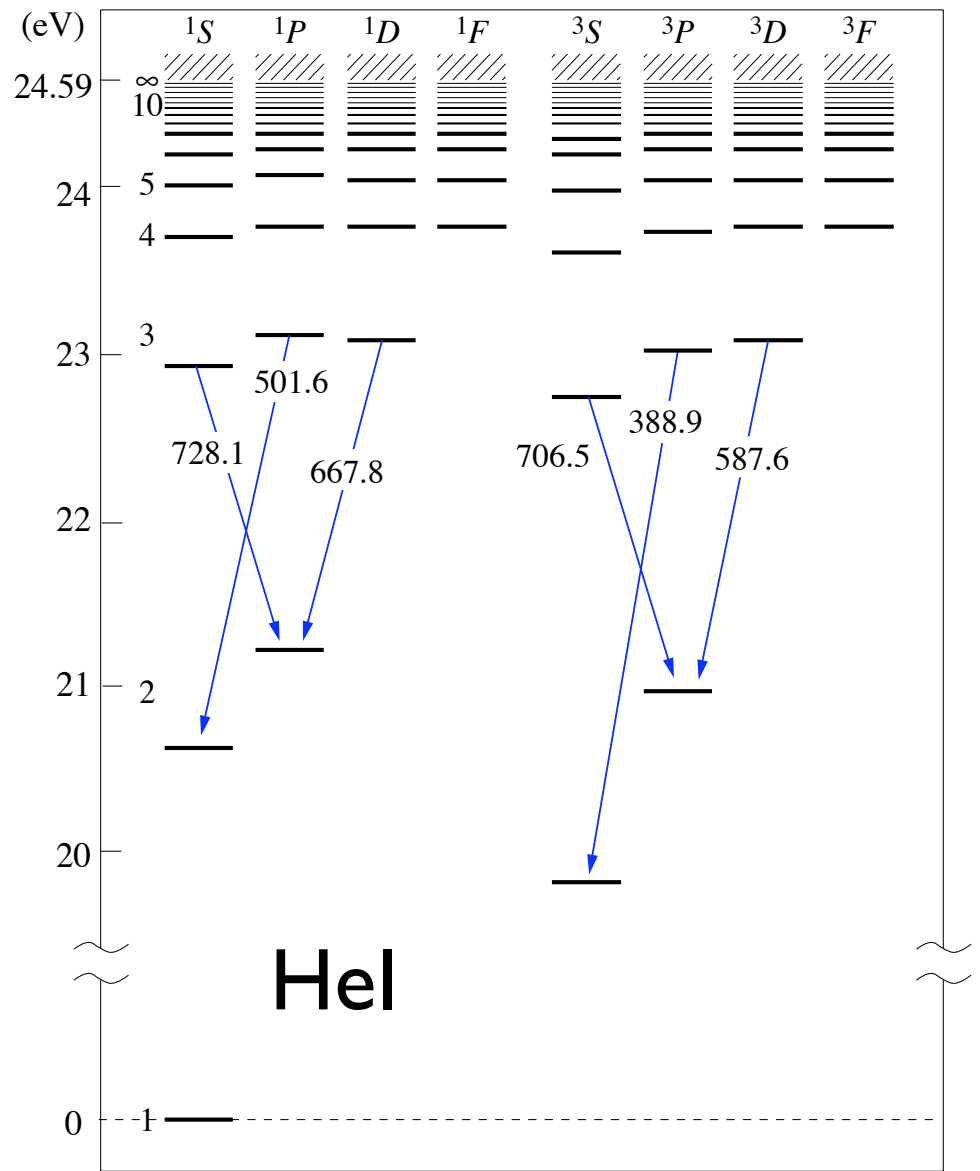
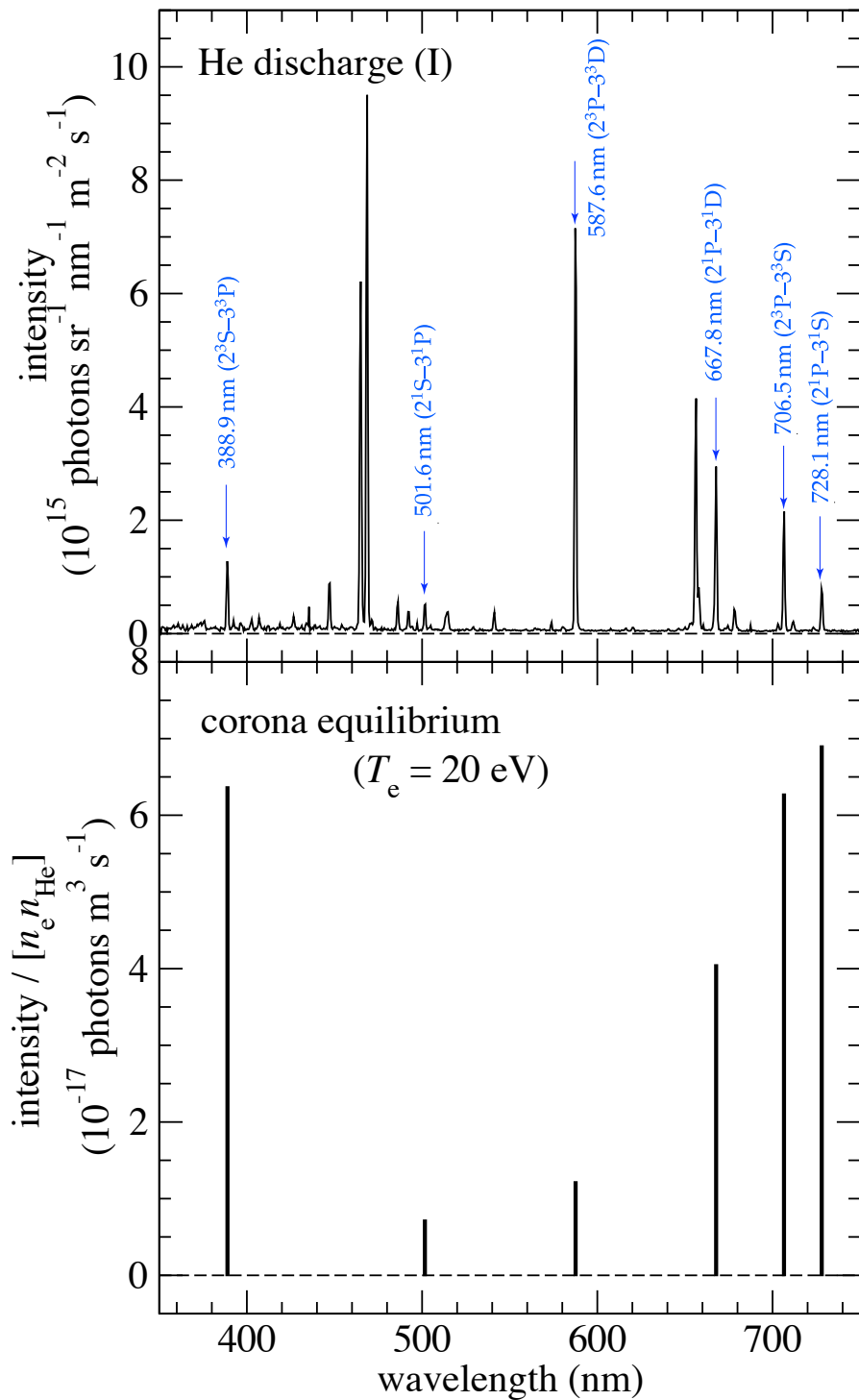
$$C(q, p)n_e n(1) \quad [\text{m}^{-3}\text{s}^{-1}]$$

rate coefficients are  
obtained from cross section

$$C(q, p) = \int_0^{\infty} \sigma_{qp}(v) f(v) v dv$$

$C(q, p)$  depends on  $T_e$

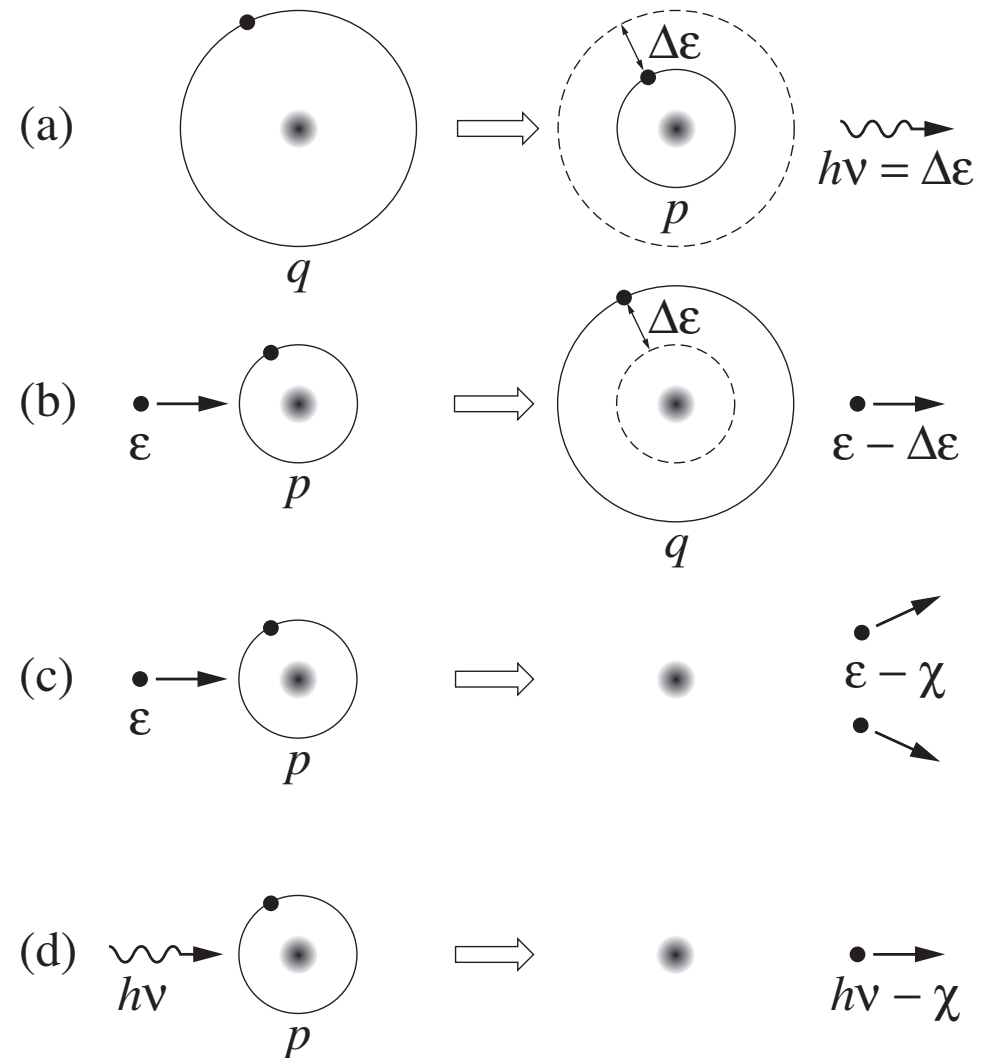
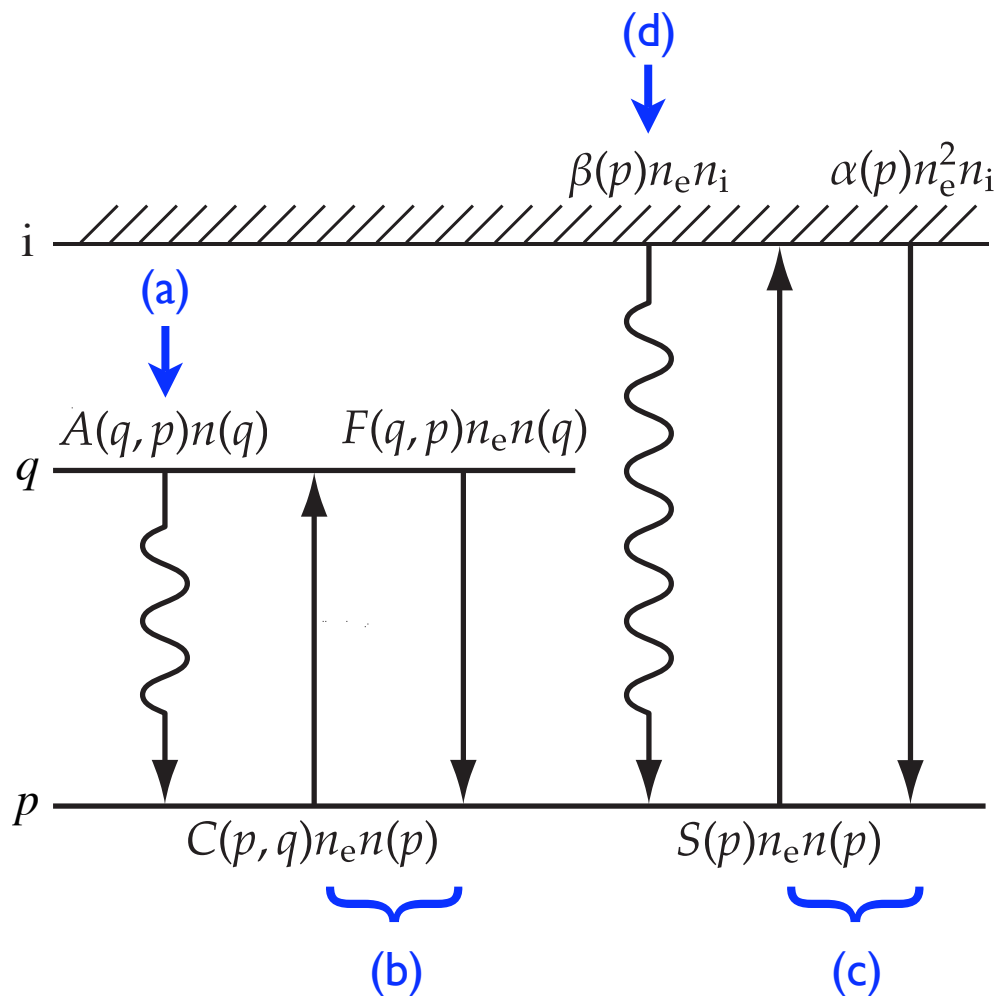




corona model does not explain observed spectrum

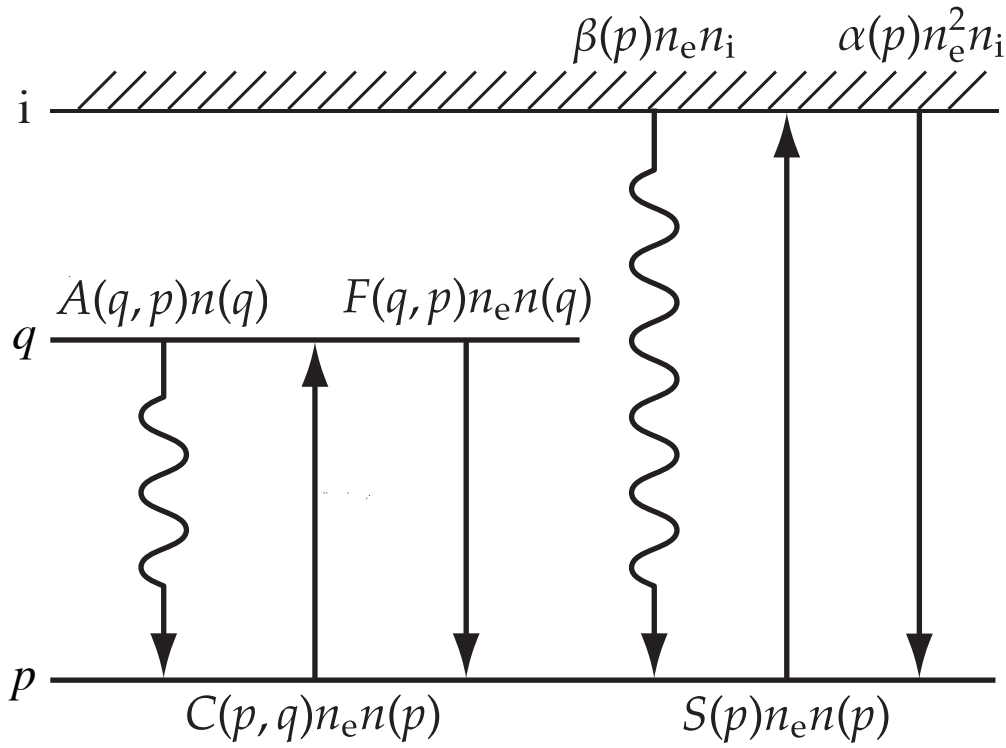
# atomic processes

various processes should be taken into consideration





# collisional-radiative model



$$\frac{d}{dt}n(p) = \Gamma_{\text{in}} - \Gamma_{\text{out}} = 0$$

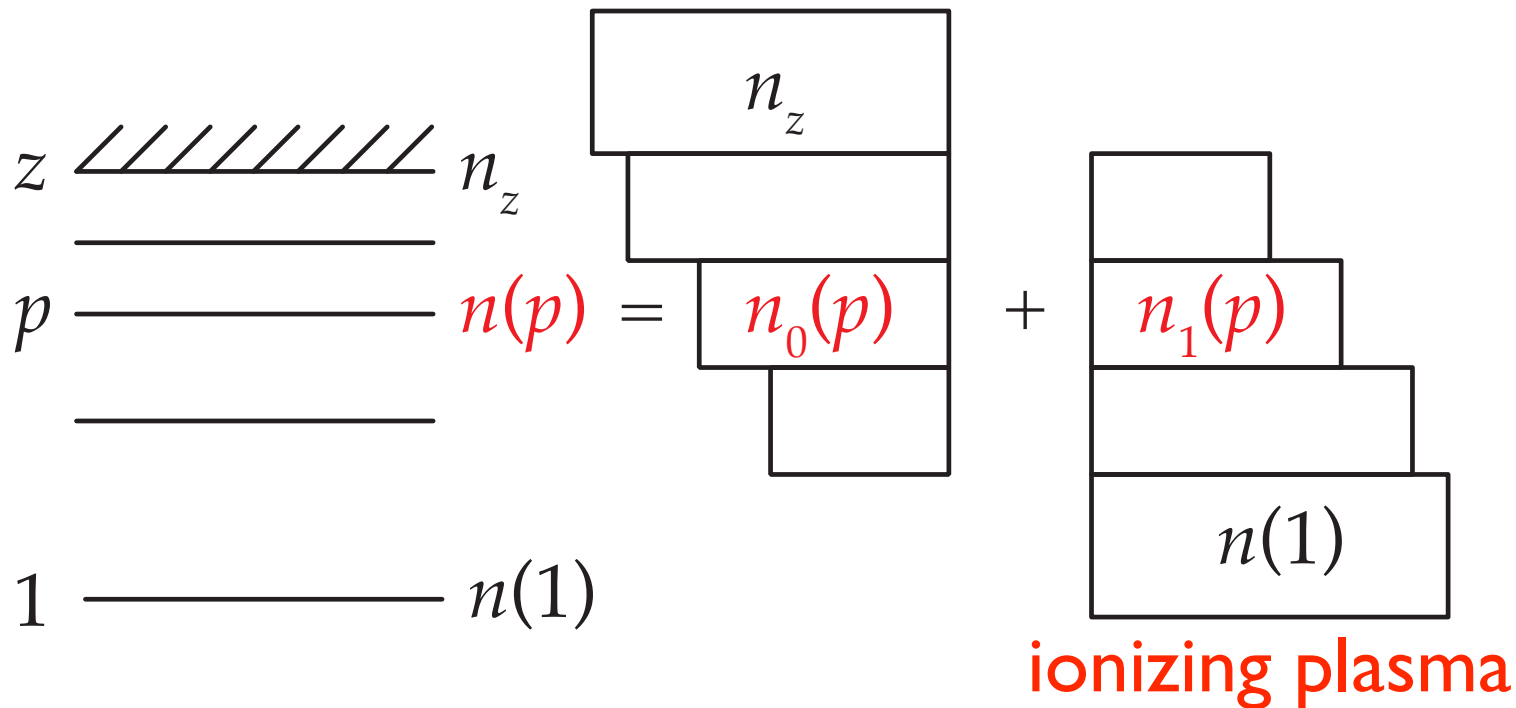
quasi steady-state approximation

$$\begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix} \begin{pmatrix} n(2) \\ \cdot \\ n(p) \\ \cdot \end{pmatrix} = \begin{pmatrix} \cdot \\ \cdot \\ \cdot \\ \cdot \end{pmatrix} n(1) + \begin{pmatrix} \cdot \\ \cdot \\ \cdot \\ \cdot \end{pmatrix} n_i$$

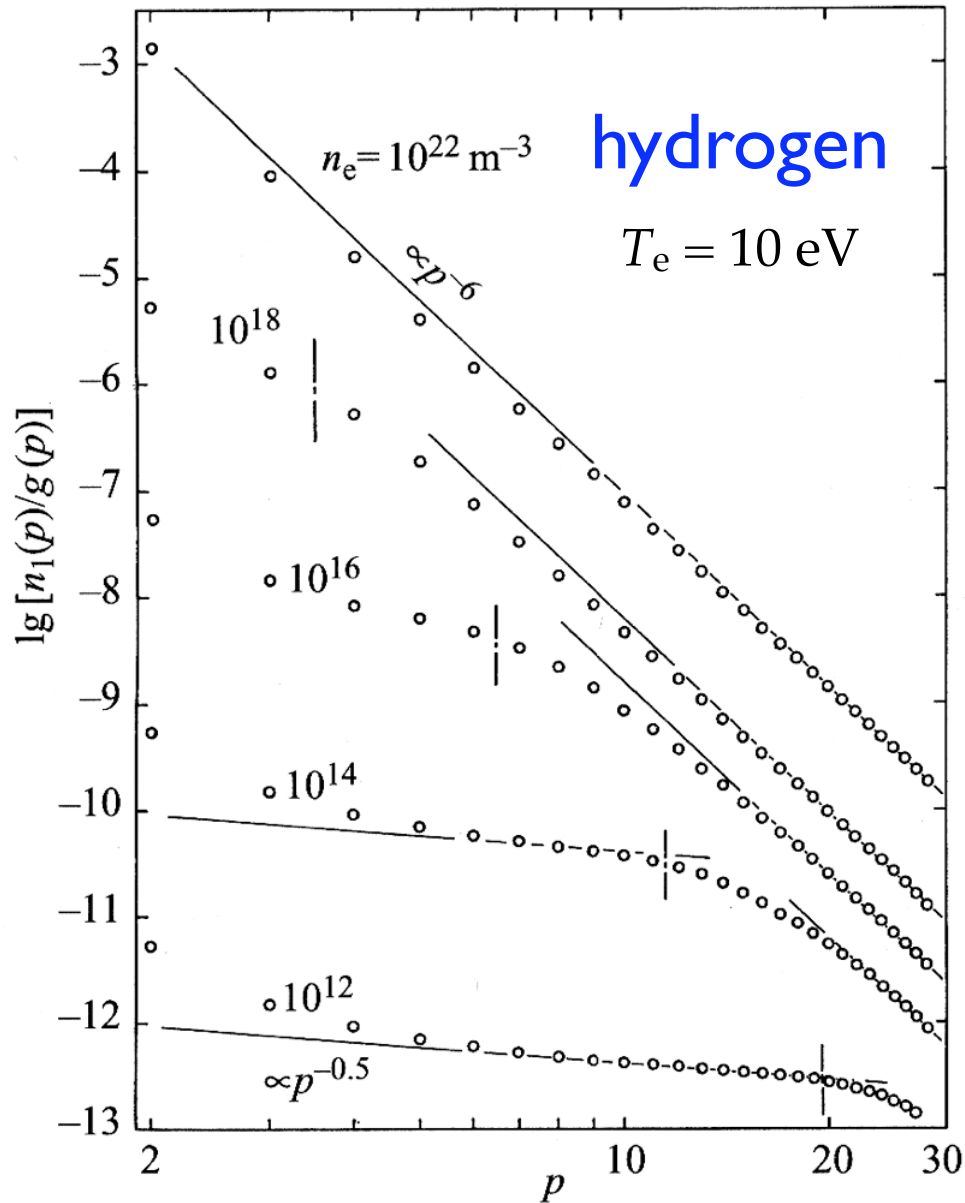
population consists of two independent components

$$\begin{aligned}n(p) &= r_0(p)n_e n_z + r_1(p)n_e n(1) \\ &= n_0(p) + n_1(p)\end{aligned}$$

recombining plasma



# ionizing plasma



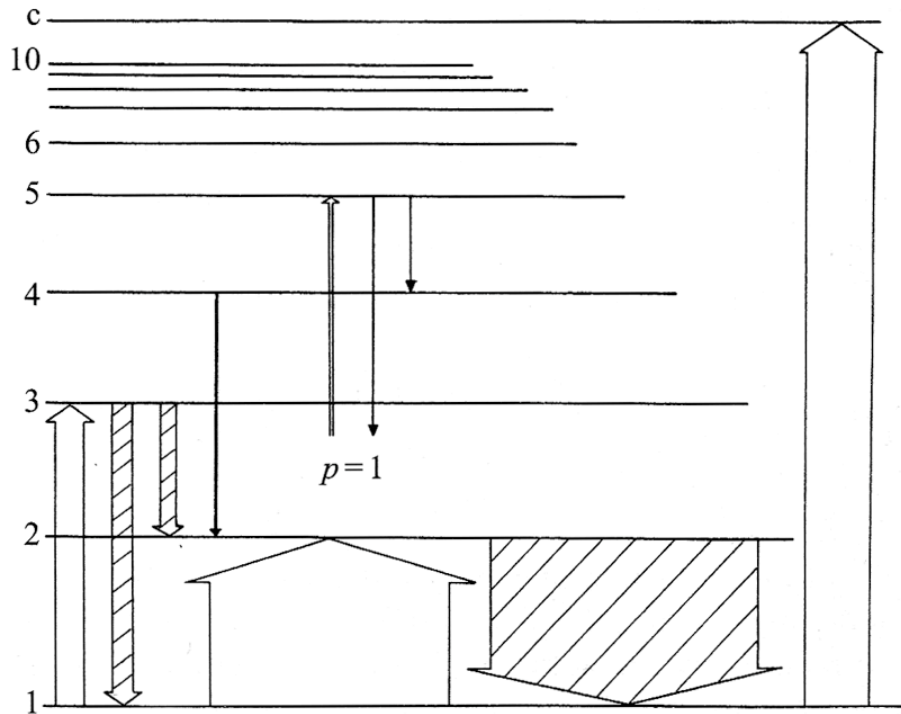
$$\frac{n(p)}{g(p)} \propto \begin{cases} p^{-0.5} & \text{for low } n_e \\ p^{-6} & \text{for high } n_e \end{cases}$$

$T_e$  dependence on  $n(p)$  distribution is small

(T. Fujimoto, *Plasma Spectroscopy*)

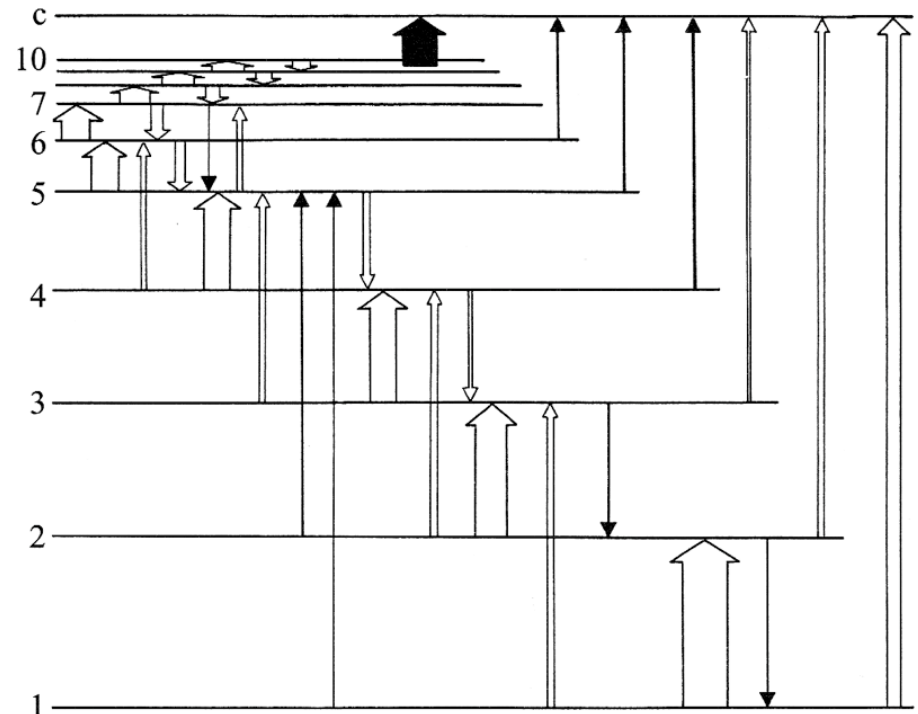
# population flows (ionizing plasma)

low  $n_e$  case



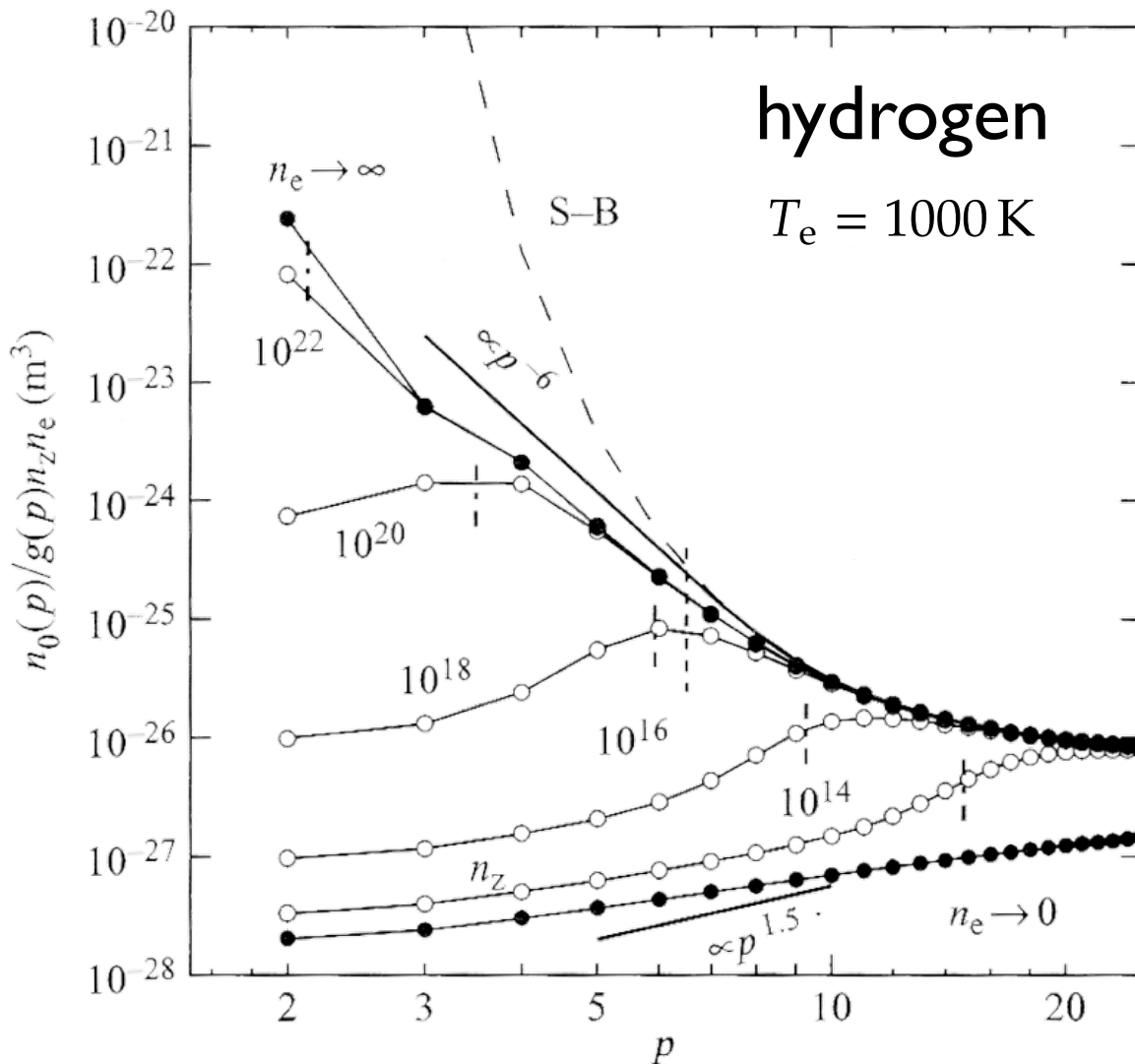
close to **corona equilibrium**

high  $n_e$  case



(T. Fujimoto, *Plasma Spectroscopy*)

# recombining plasma



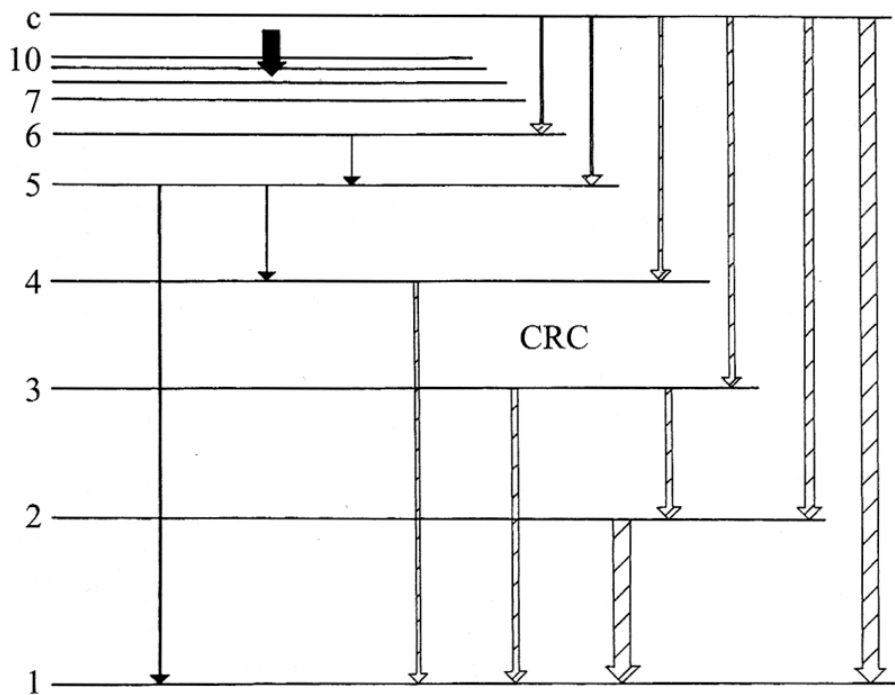
when  $n_e$  is high,  $n(p)$  is expressed with **Saha-Boltzmann equation**

$$n(p) = \frac{g(p)}{2g_z(1)} \left( \frac{h^2}{2\pi m k T_e} \right)^{3/2} \exp \left[ \frac{\chi(p)}{k T_e} \right] n_e n_z(1)$$

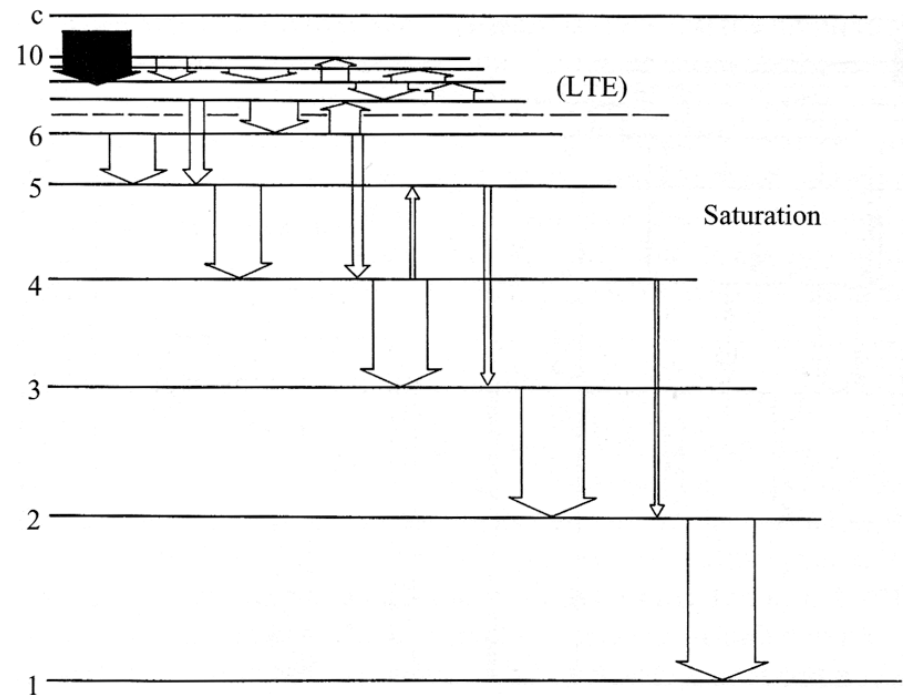
(T. Fujimoto, *Plasma Spectroscopy*)

# population flows (recombining plasma)

low  $n_e$  case

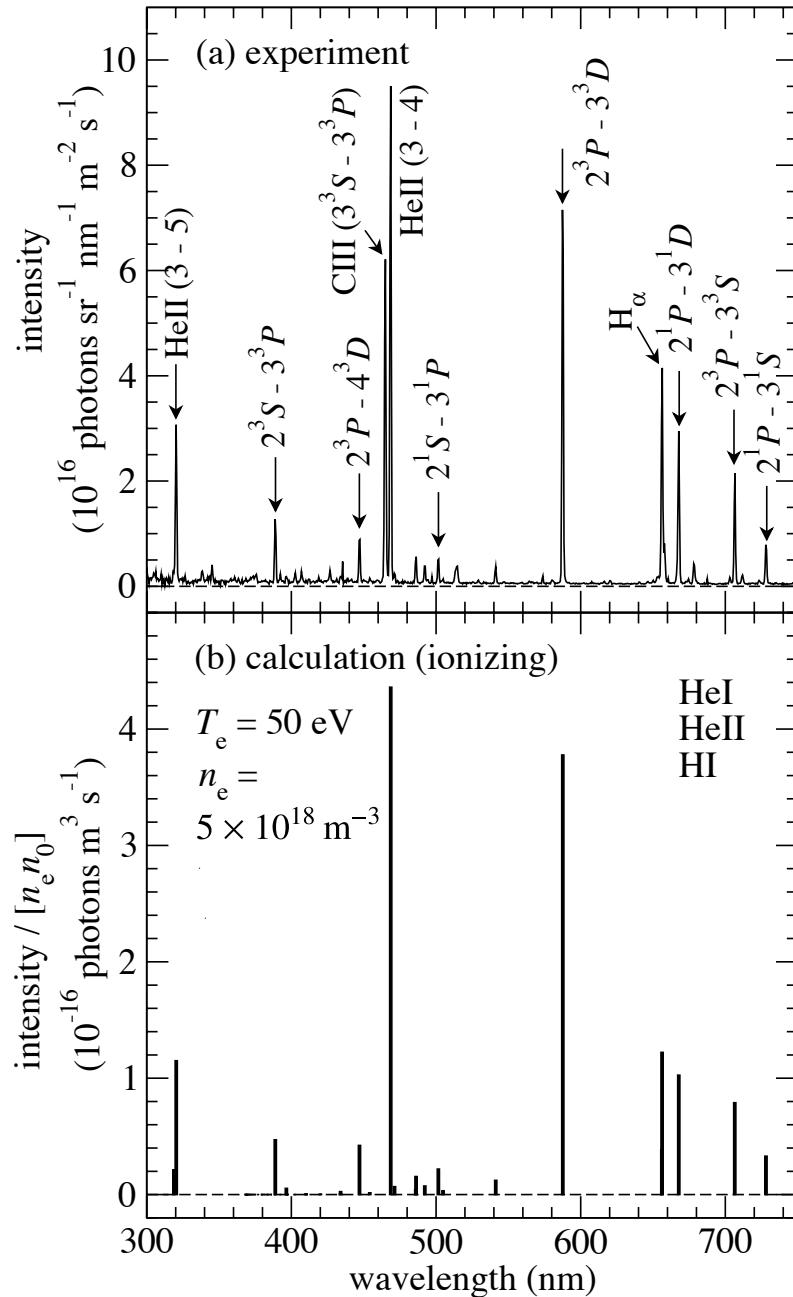


high  $n_e$  case

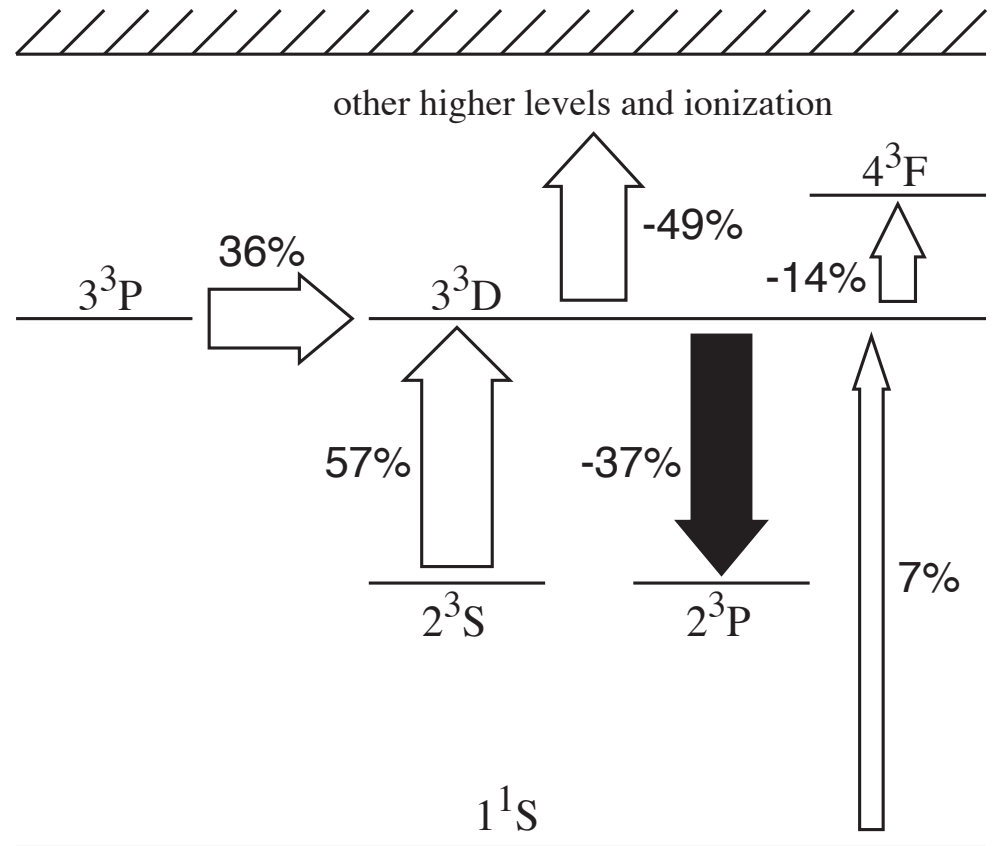


(T. Fujimoto, *Plasma Spectroscopy*)

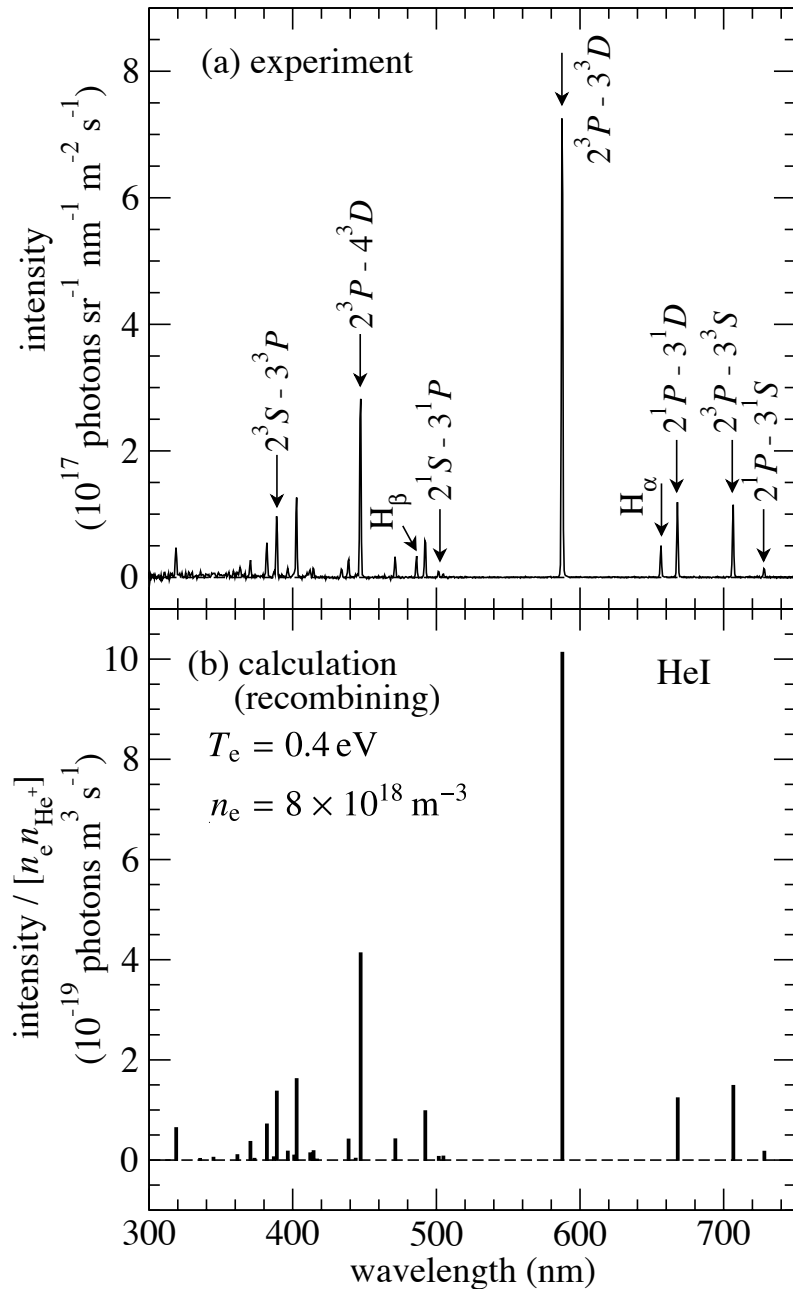
# ionizing plasma (I)



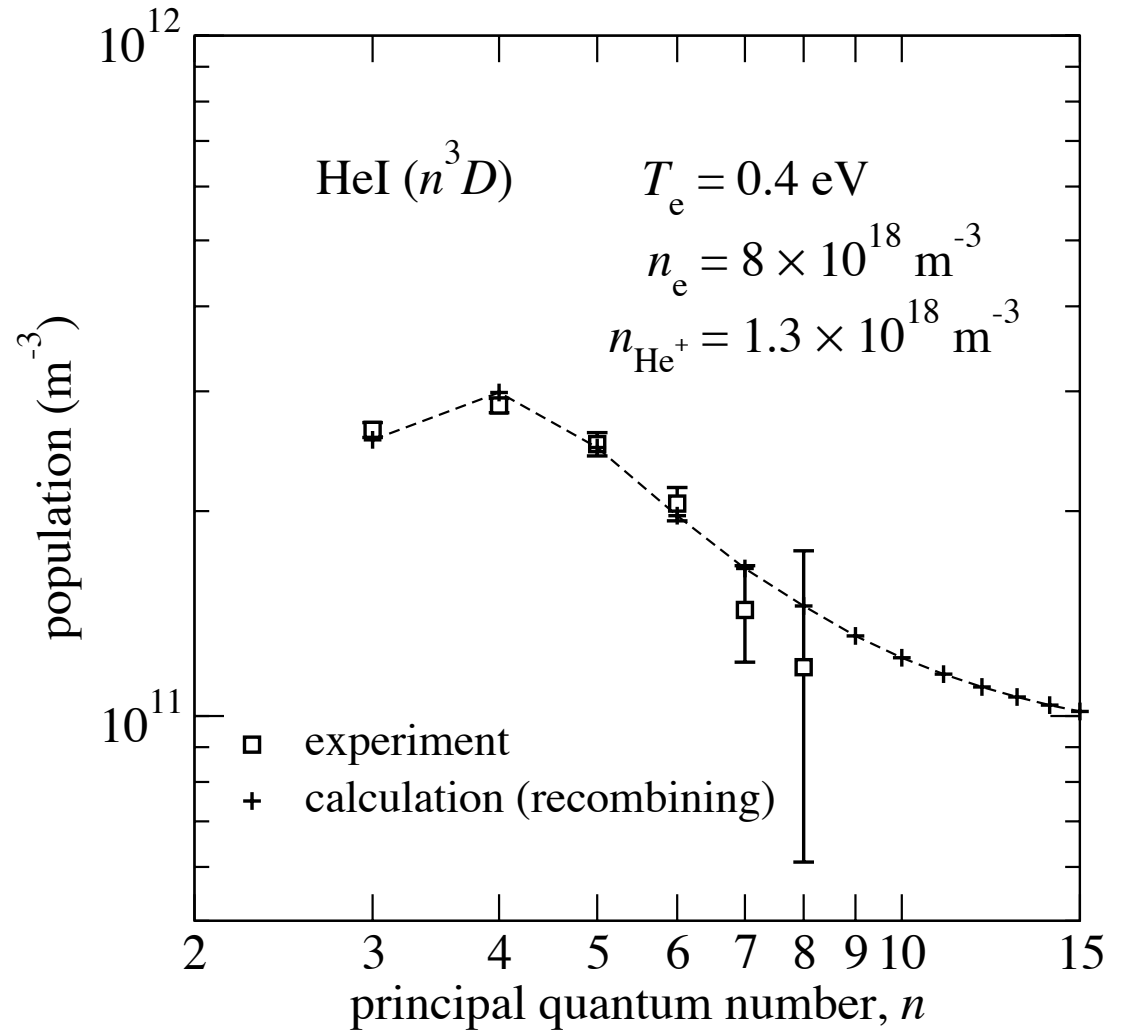
$$n(p) = r_1(p)n_e n(1)$$



# recombining plasma (III)

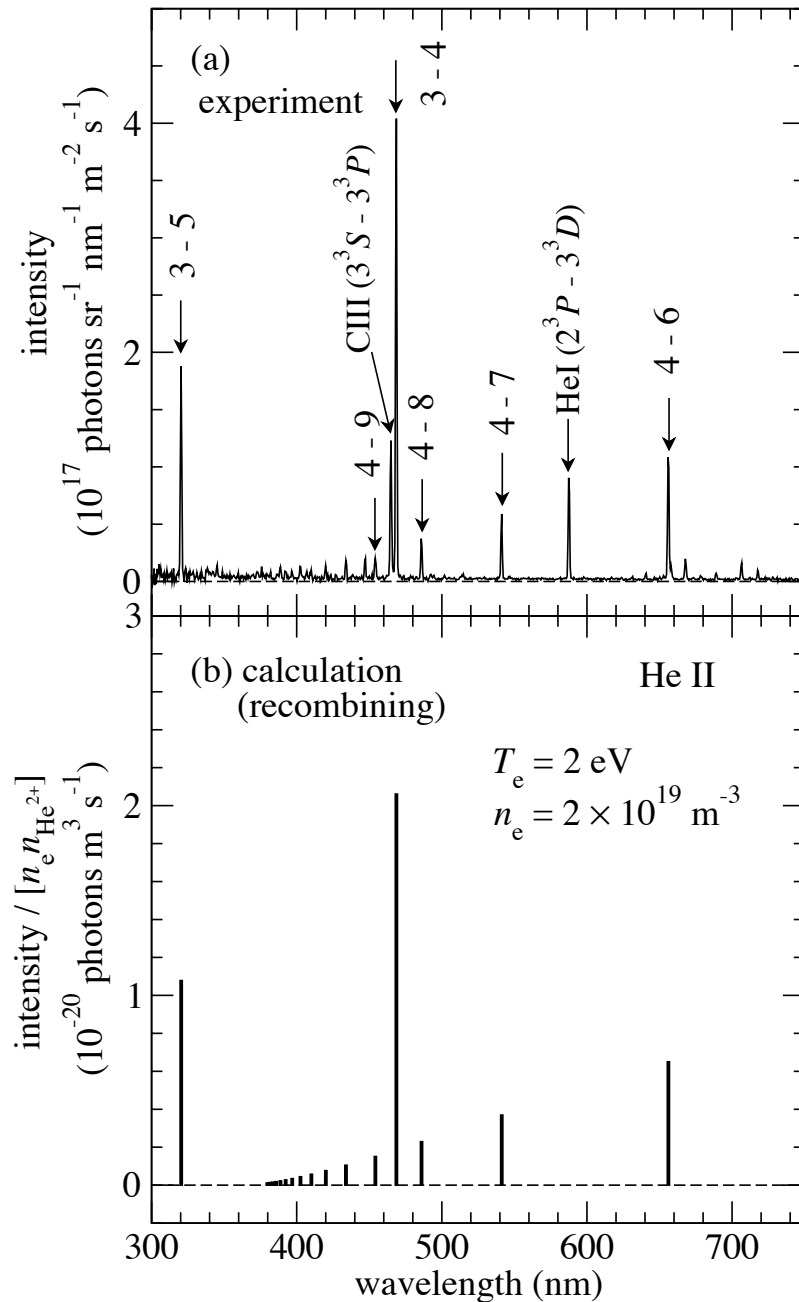


$$n(p) = r_0(p)n_en_z$$





# recombining plasma (II)



- recombining plasma of HeII appears earlier than HeI
- derived  $T_e$  is higher

**effect of external fields**

# splitting, shift, broadening, etc

- Zeeman effect — magnetic field
- Stark effect — electric field
- Stark broadening — electric micro fields

# Zeeman effect

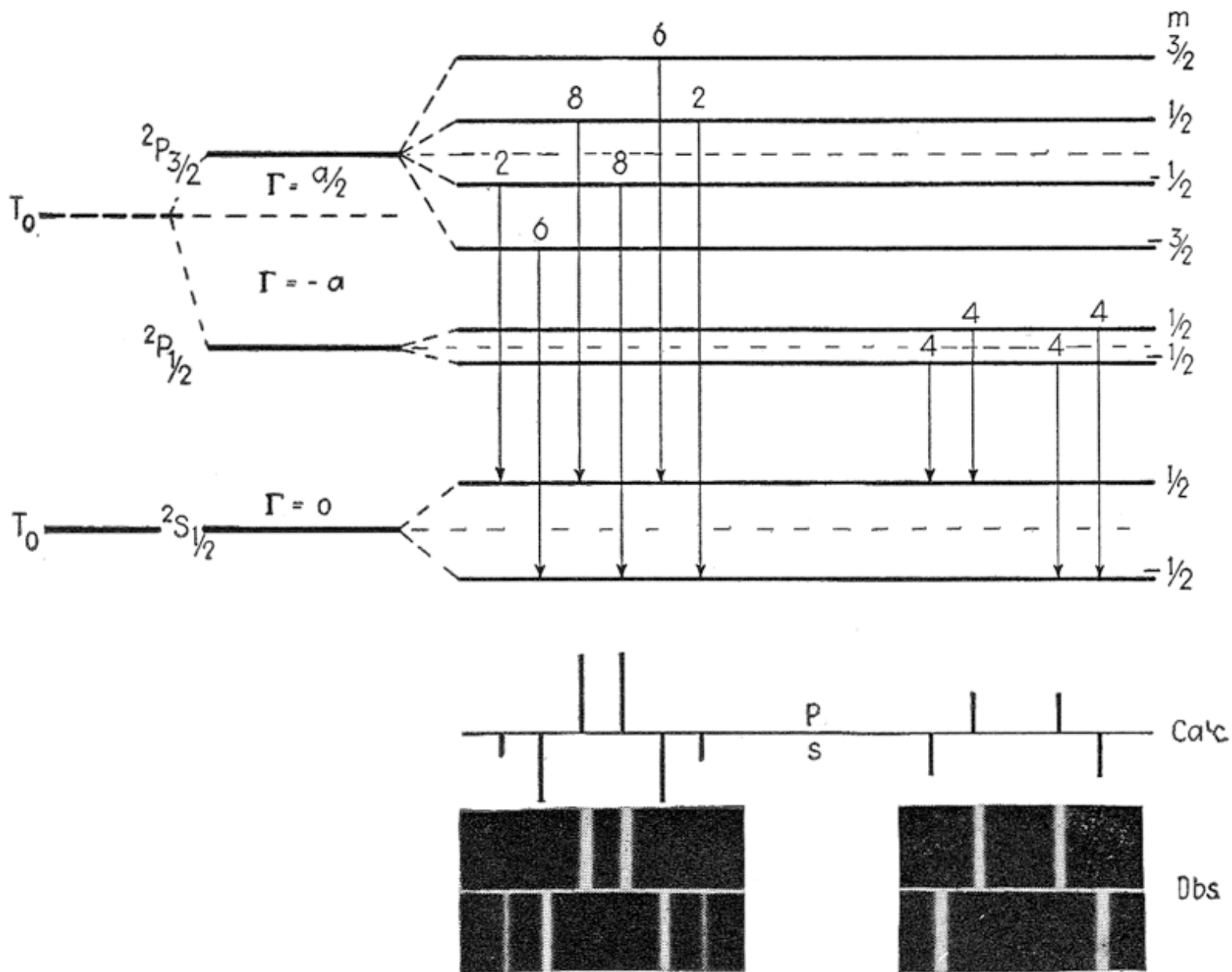


FIG. 10.8.—Zeeman effect of a principal-series doublet.

(H. E. White, *Introduction to Atomic Spectra*)

# polarization

$\Delta M = \pm 1 \rightarrow$  polarization perpendicular to  $B$  ( $\sigma$ -polarization)  
 $\Delta M = 0 \rightarrow$  polarization parallel to  $B$  ( $\pi$ -polarization)

circularly polarized light ( $\sigma$ -polarization)

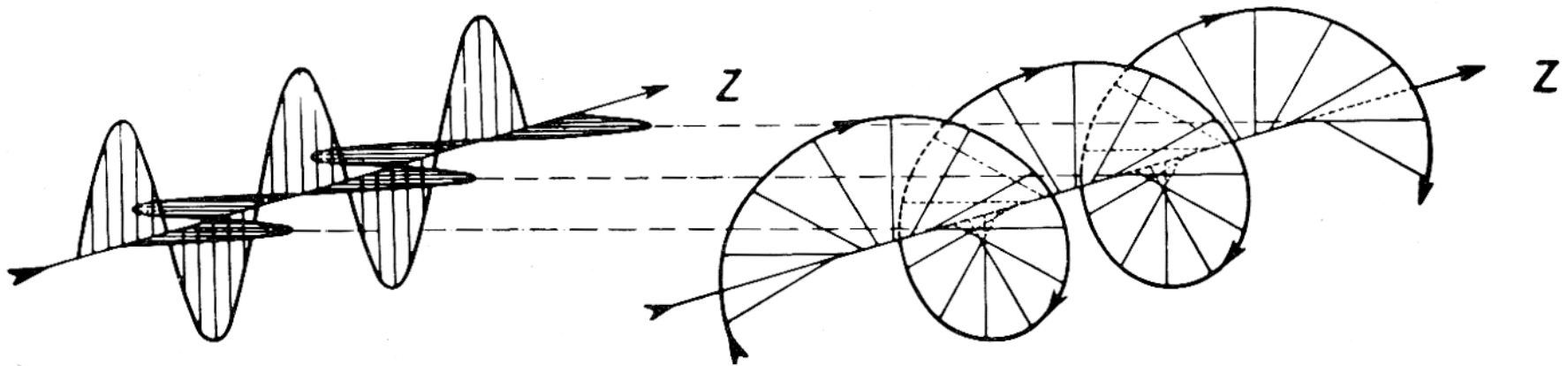


Fig. 10.2. Superposition of two linear-polarized phase-shifted waves (left-hand figure) into a circularly-polarized wave (right-hand figure). The sense of rotation is defined in the text.

(P. H. Heckmann, *Introduction to the Spectroscopy of Atoms*)

$\pi$ -component is not seen for  $// B$  observation

$\pi$  and  $\sigma$  have the same intensity for  $\perp B$  observation

# perturbation theory (eigensystems)

$$H\psi = E\psi$$

$$H = H^0 + V$$

$$= \begin{pmatrix} E_1^0 + V_{11} & V_{12} & \cdots & V_{1n} \\ V_{21} & E_2^0 + V_{22} & \cdots & V_{2n} \\ \vdots & \vdots & \ddots & \vdots \\ V_{n1} & V_{n2} & \cdots & E_n^0 + V_{nn} \end{pmatrix}$$

# perturbation of magnetic field

$|LSJM\rangle$  is taken as base function

$$\langle JM|V|J'M'\rangle$$

$$= -\mu_B B \langle JM|(g_L L_z + g_S S_z)|J'M'\rangle$$

$$= -\mu_B B \sum_{M_L M_S} \langle JM|M_L M_S\rangle \langle J'M|M_L M_S\rangle (g_L M_L + g_S M_S)$$

$$= -\mu_B B \sqrt{2J+1} \sqrt{2J'+1}$$

$$\sum_{M_L M_S} \begin{pmatrix} 1 & 1 & J \\ M_L & M_S & -M \end{pmatrix} \begin{pmatrix} 1 & 1 & J' \\ M_L & M_S & -M \end{pmatrix} (g_L M_L + g_S M_S)$$

$$(M = M_L + M_S)$$

$$1s2p\ ^3P_{0,1,2} \quad \langle JM|V|J'M'\rangle = -\mu_B B \langle JM|(g_L L_z + g_S S_z)|J'M'\rangle$$

	$J$	0	1		2					
$J'$	$M$	0	-1	0	1	-2	-1	0	1	2
0	0		-0.8165							
	-1		-1.5000			-0.5000				
1	0	-0.8165						-0.5774		
	1				1.5000			-0.5000		
	-2				-3.0000					
	-1		-0.5000			-1.5000				
2	0		-0.5774							
	1				-0.5000			1.5000		
	2								3.0000	





perturbed eigenstate is expressed with base functions

$$|\psi_m\rangle = \sum_m C_i^m |\psi_i^0\rangle$$

transition probability is written as

$$A_{\psi_m \rightarrow \varphi_n}^q \propto |\langle \varphi_n | D_q | \psi_m \rangle|^2 = \left| \sum_{i,j} C_j^n C_i^m \langle \varphi_j^0 | D_q | \psi_i^0 \rangle \right|^2$$

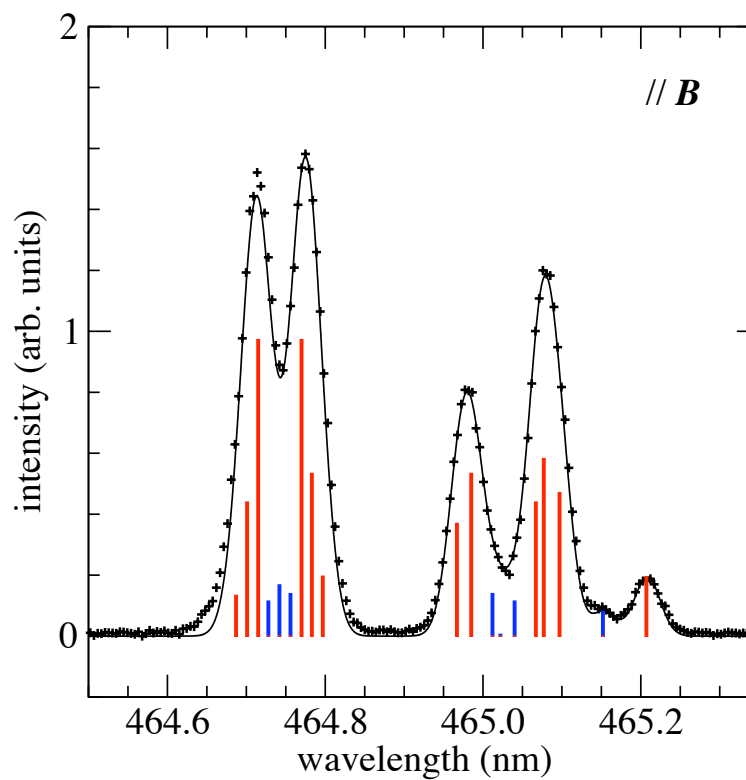
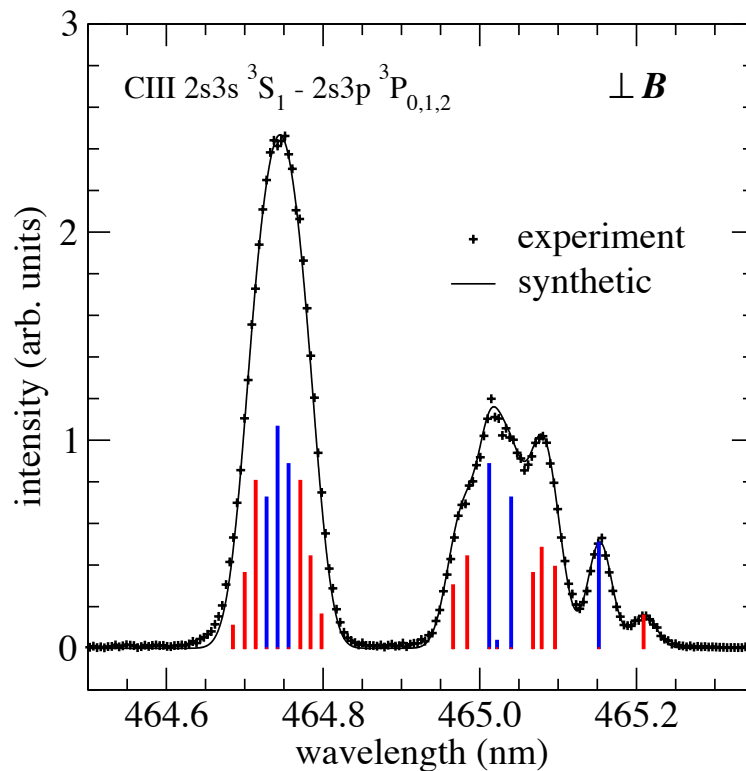
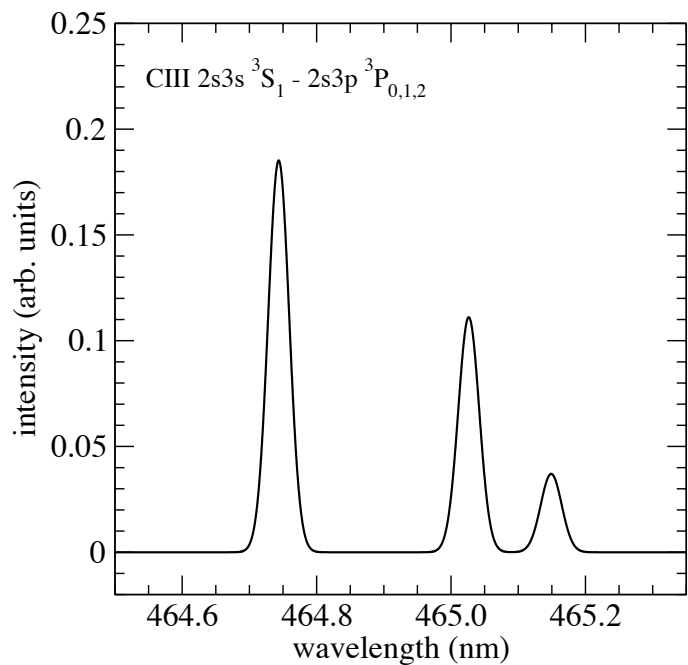
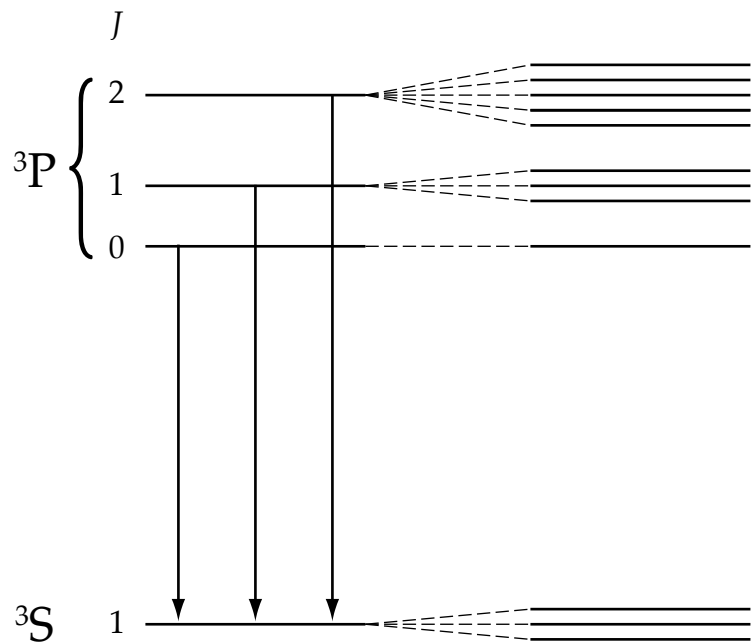
each term is calculated as

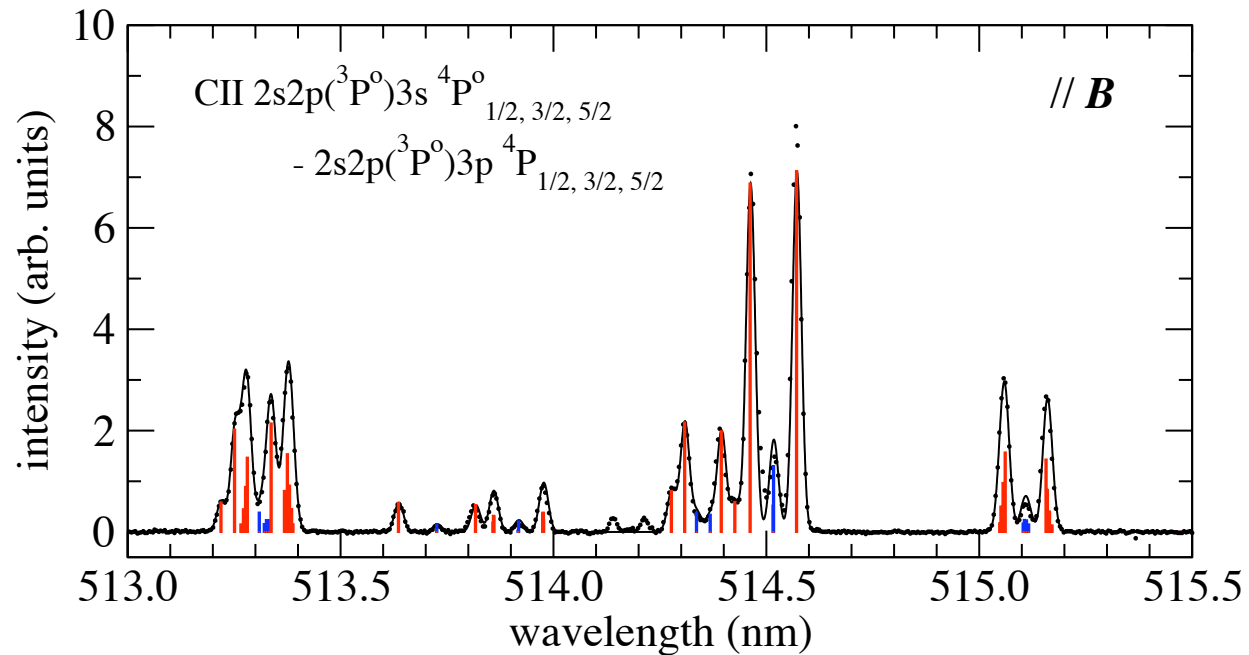
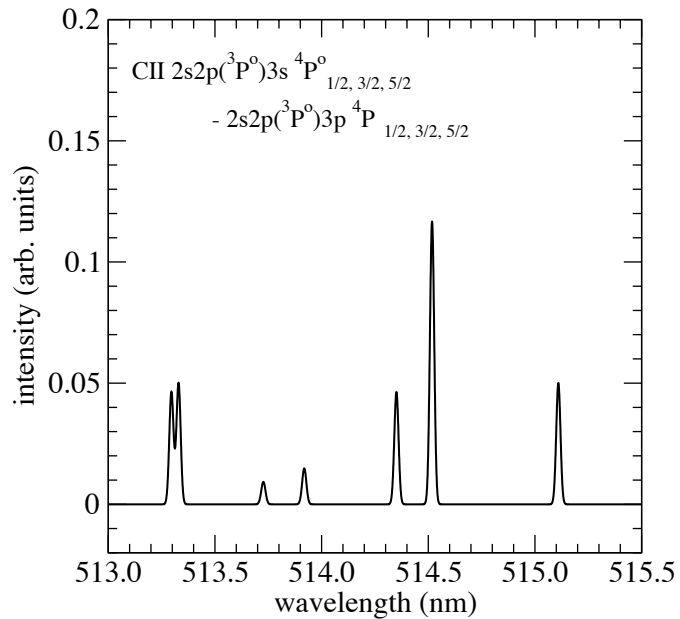
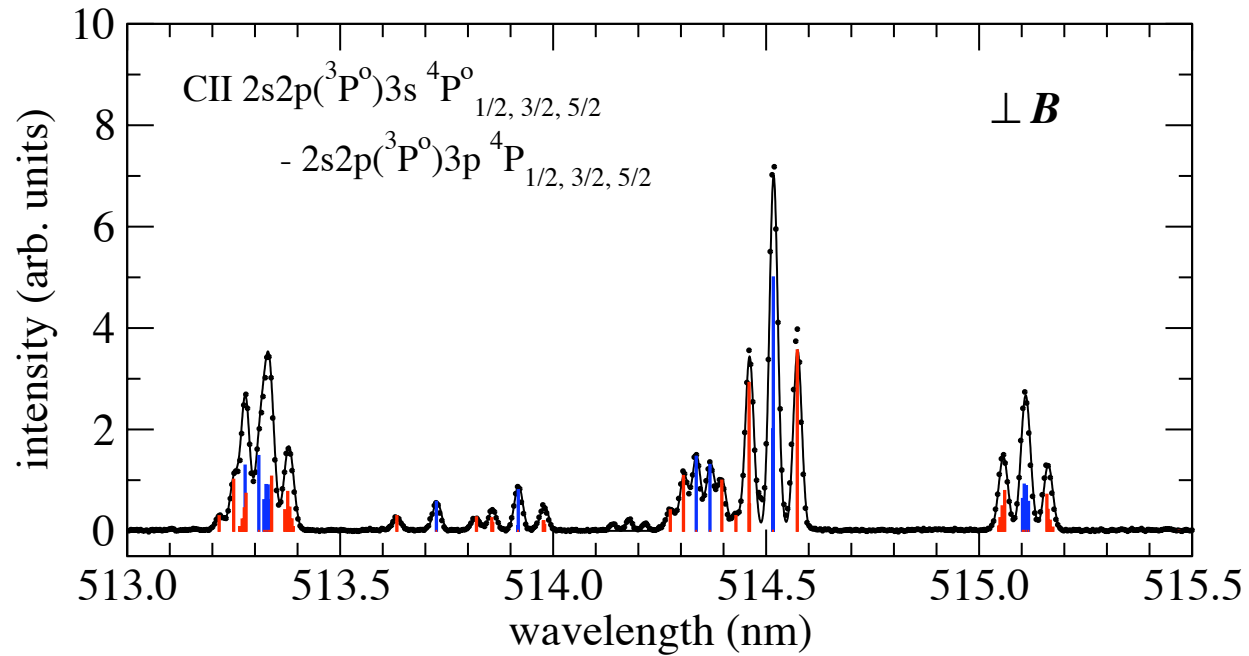
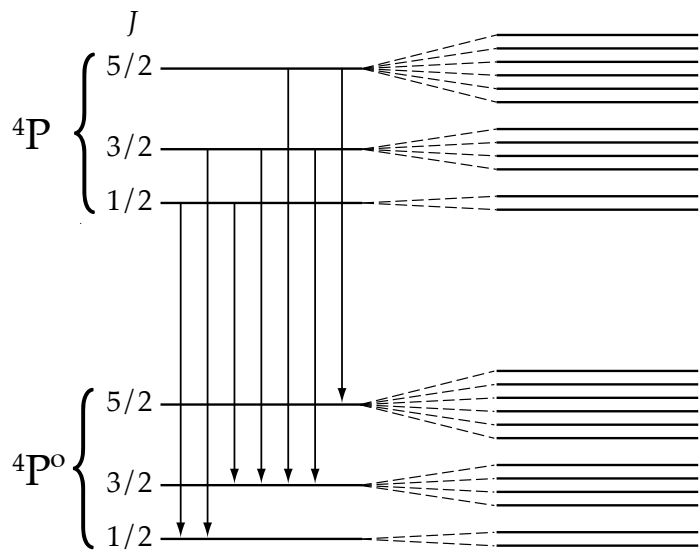
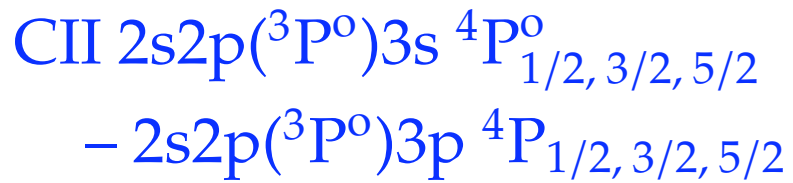
$$|\psi_i^0\rangle = |L_1 S_1 J_1 M_1\rangle$$

$$|\varphi_j^0\rangle = |L_2 S_2 J_2 M_2\rangle$$

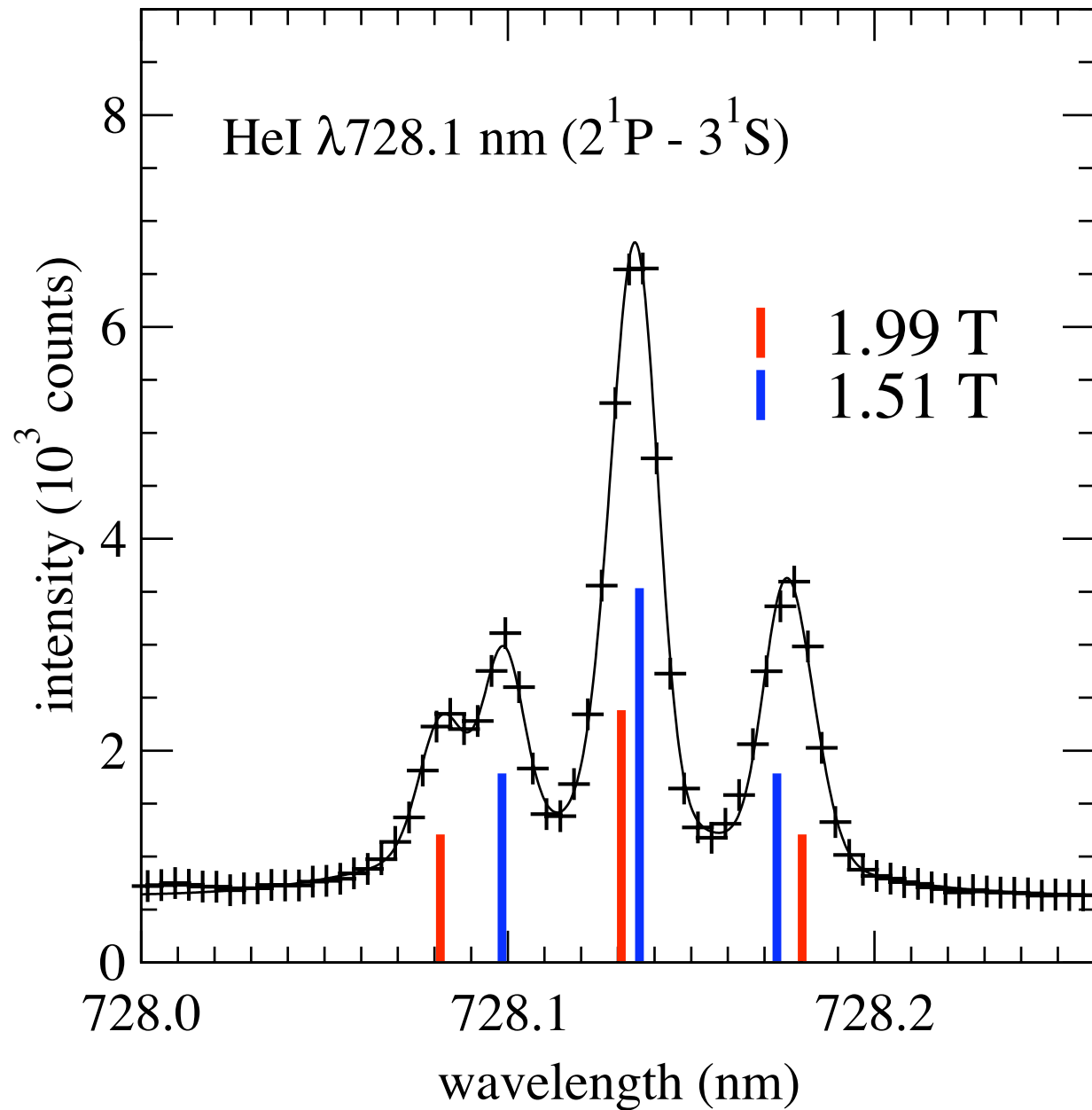
$$\langle \varphi_j^0 | D_q | \psi_i^0 \rangle = (-1)^{(J_1 - M_1) + (S + 1 + L_1 + J_2)} \sqrt{(2J_1 + 1)(2J_2 + 1)} \\ \left( \begin{array}{ccc} J_1 & 1 & J_2 \\ -M_1 & q & M_2 \end{array} \right) \left\{ \begin{array}{ccc} L_1 & J_1 & S \\ J_2 & L_2 & 1 \end{array} \right\} \langle L_1 || D || L_2 \rangle$$

# CIII $2s3s\ ^3S_1 - 2s3p\ ^3P_{0,1,2}$

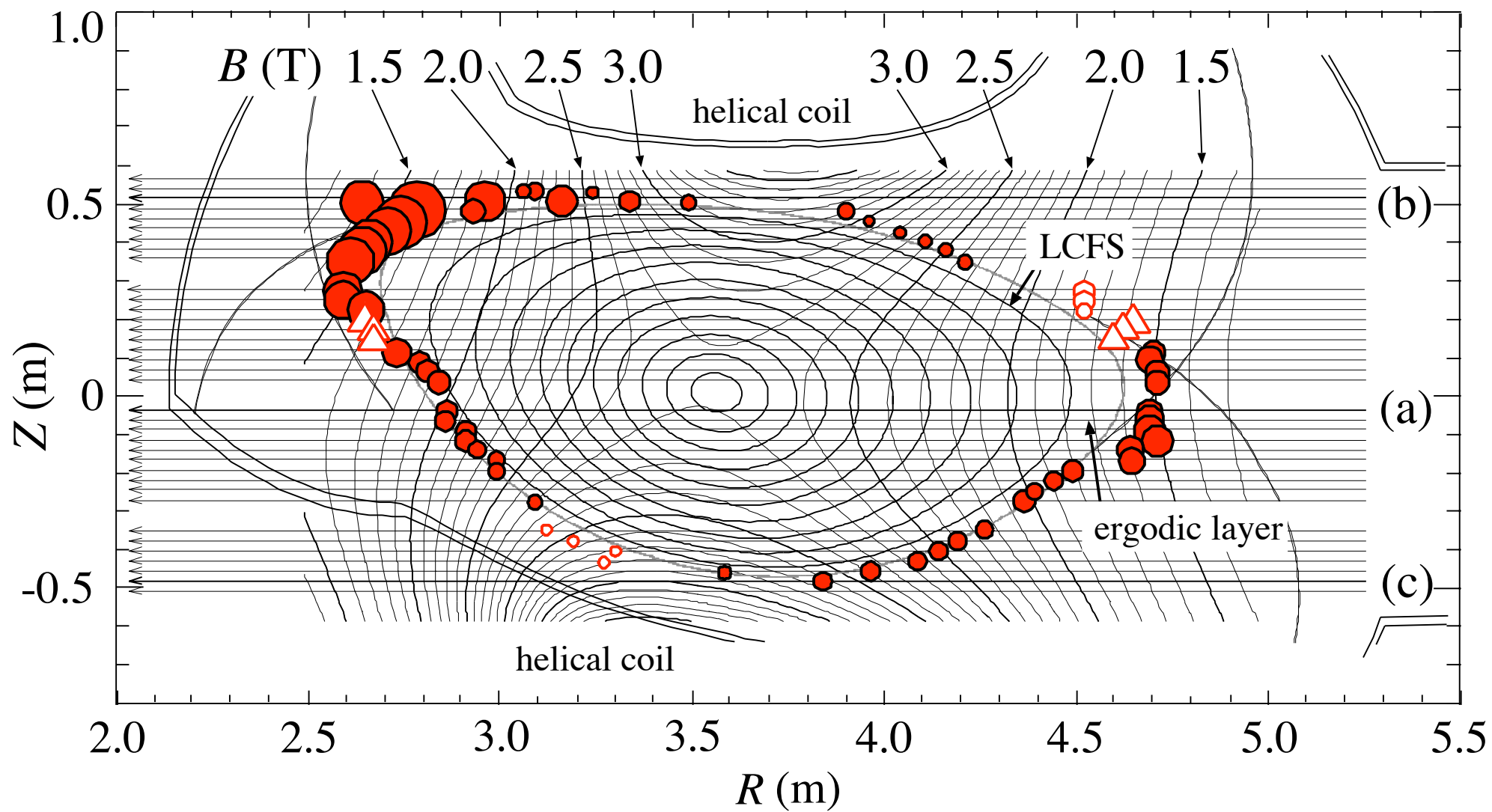




# application to fusion plasma



(M. Goto et al., Phys. Rev. E 65, 026401 (2002))



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